Fractionation of ²³⁰Th, ²³¹Pa, and ¹⁰Be induced by particle size and composition within **an opal-rich sediment of the Atlantic Southern Ocean**

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Abstract

This study centers on the question: How sensitive are ²³¹Pa^{230}Th and ¹⁰Be^{230}Th to sediment composition and redistribution? The natural radionuclides ^{231}Pa , ^{230}Th and ^{10}Be recorded in deep sea sediments are tracers for water mass advection and particle fluxes. We investigate the influence of oceanic particle composition on the element adsorption in order to improve our understanding of sedimentary isotope records. We present new data on particle size specific 231 Pa and 10 Be concentrations. An additional separation step, based on settling velocities, led to the isolation of a very opal-rich phase. We find that opal-rich particles contain the highest ²³¹Pa and ¹⁰Be concentrations, and higher ²³¹Pa^{$/230$}Th and ¹⁰Be^{$/230$}Th isotope ratios than opal-poor particles. The fractionation relative to 230 Th induced by the adsorption to opal-rich particles is more pronounced for 231 Pa than for 10 Be. We conclude that bulk $^{231}Pa^{230}Th$ in Southern Ocean sediments is most suitable as a proxy for past opal fluxes. The comparison between two neighboring cores with rapid and slow accumulation rates reveals that these isotope ratios are not influenced significantly by the intensity of sediment focusing at these two study sites. However, a simulation shows that particle sorting by selective removal of sediment (winnowing) could change the isotope ratios. Consequently, $^{231}Pa^{230}Th$ should not be used as paleocirculation proxy in cases where a strong loss of opalrich material due to bottom currents occurred.

1 **1. INTRODUCTION**

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1.1 Natural radionuclides and their application as proxies

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5 6 7 8 9 10 11 12 The natural radionuclides ^{231}Pa , ^{230}Th , and ^{10}Be are powerful tracers, used to identify and quantify the transport of particles and of water masses in the ocean. The common property of Th, Pa, and Be in seawater is their particle reactivity, which makes them suitable for tracing past particle fluxes. They adsorb onto particles, sink out of the water column (scavenging), and are buried in the sediment. Pa and Be have a lower reactivity and consequently longer residence time than Th. This differing solubility creates disequilibria, and finally leads to isotopic ratios that differ from their production ratio, depending on their residence times in water column.

13 14 15 16 17 18 19 20 21 22 23 24 The production ratios of U-series nuclides are well defined. 230 Th (half-life 75,380) years) is produced by its progenitor ²³⁴U in seawater (~2910 dpm m^{-3} at 35 permil salinity, Robinson et al., 2004) at a constant rate of 0.0267 dpm m^{-3} a^{-1} (François et al., 2004). Due to its strong particle reactivity (Langmuir and Herman, 1980), 230 Th is adsorbed to colloids or suspended particles shortly after its production in seawater (Moore and Hunter, 1985; Moore and Millward, 1988; Niven and Moore, 1993). The scavenging process for thorium is very efficient so that the rain rate of particulate 230 Th to the sea floor is nearly independent from the intensity of particle fluxes, and thus the vertical 230 Th flux is spatially and temporally constant and equals its production in the overlying water column (within an uncertainty of ± 30 %, Henderson et al., 1999; Scholten et al., 2001). Therefore, the residence time of ²³⁰Th in the water column with respect to scavenging is very short (10-40 years, Anderson et al., 1983a, 1983b). Due to its nearly constant vertical flux, the 230 Th concentration in sediments is 25 26 inversely related to the intensity of the vertical particle flux (Krishnaswami, 1976; Bacon, 1984; Suman and Bacon, 1989).

27 28 29 30 31 32 33 34 35 36 37 38 The source of ²³¹Pa (half-life 32,760 years) in seawater is the radioactive decay of ²³⁵U at a constant rate of 0.0025 dpm $m^{-3} a^{-1}$. The activity ratio of $2^{31} Pa^{230}$ Th at the time of production from 235 U and 234 U is 0.093. However, protactinium is less reactive and has a longer scavenging residence time than thorium (10² years, Anderson et al., 1983a; 1983b). Therefore, $^{231}Pa^{230}Th$ ratios may differ from the production ratio in response to the effects that lead to their removal and transport. In contrast to 230 Th, the distribution of dissolved 231 Pa is largely influenced by lateral mixing and advection of water masses that leads to its further transport and deposition in regions with enhanced particle fluxes (boundary scavenging, Anderson et al., 1983a, 1983b). Due to its sensitivity to the hydrography, the ²³¹Pa α ²³⁰Th ratio in sediments has been interpreted as a record of deep water ventilation in the Atlantic (Yu et al., 1996; McManus et al., 2004; Gherardi et al., 2009; Guihou et al., 2010; Negre et al., 2010; Guihou et al., in press).

39 40 41 42 43 44 45 46 47 48 49 50 As the ²³¹Pa α ²³⁰Th ratio is positively correlated to the intensity of particle fluxes (Yang et al., 1986), it has also been applied as an export productivity proxy (Kumar et al., 1993; Pichat et al., 2004). However, in the Southern Ocean the scavenging of 231 Pa has been found to be mainly controlled by opal flux (Rutgers van der Loeff and Berger, 1993; Walter et al., 1997). Therefore, in most recent studies located in high latitudes, the application of $^{231}Pa^{230}Th$ has been restricted to reconstructing past changes of opal fluxes (Bradtmiller et al., 2007, 2009; Anderson et al., 2009), rather than total mass fluxes. It is still a matter of debate if the sensitivity of $^{231}Pa^{230}Th$ ratios to particle composition affects their application as paleocirculation tracers, particularly in the North Atlantic. However, modeling results point to a significant role of particle composition for $^{231}Pa^{230}Th$ records (Siddall et al. 2005), and the general lack of larger sets of observational data in combination with useful particle characterizations is hampering the reliable application of the proxy (Burke et al., 2011).

70 **1.2 The effect of sediment redistribution on nuclide ratios**

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72 73 74 75 76 It has been questioned to what extent the sedimentary $^{231}Pa^{230}Th$ and $^{10}Be^{230}Th$ records are determined by ventilation versus particle composition. Indeed, recent publications (Keigwin and Boyle, 2008; Scholten et al., 2008; Gil et al., 2009; Lippold et al., 2009) claimed that the distribution of $^{231}Pa/^{230}Th$ ratios is controlled at least partly by the abundance of diatoms and not by ventilation. Various studies have shown that thorium preferentially

77 78 79 80 81 82 83 84 adsorb on lithogenics and carbonates, beryllium on lithogenics and opal, whereas protactinium strongly prefers opal (Lao et al., 1993; Chase et al., 2002; Guo et al., 2002; Chase et al., 2003; Geibert and Usbeck, 2004). The disparity in adsorption behavior and past changes in particle composition may be disadvantageous for particle flux studies (Frank et al., 2000; Gil et al., 2009; Lippold et al., 2009), as it is one additional controlling factor on the tracer flux to the sediment, which may overprint the effect of other controlling factors (such as total particle flux or water mass advection) and complicate the interpretation of the sedimentary tracer records (Burke et al., 2011).

85 86 87 88 89 90 91 92 93 94 95 96 97 The opal-rich sediments of the Southern Ocean act as a sink for 231 Pa and 10 Be as a result from hydrography, scavenging residence times, and chemical composition of oceanic particles. Opal-rich sediments appear to be particularly susceptible to resuspension, which makes redistribution by strong abyssal bottom currents particularly important in the opal-belt of the Antarctic Circumpolar Current (Geibert et al., 2005). ²³¹Pa α^{230} Th and ¹⁰Be α^{230} Th ratios have been commonly assumed to be insensitive to sediment redistribution, an assumption which remains to be proven. Focusing and winnowing of sediments occur under the influence of variable abyssal current velocities. Therefore, laterally transported sediment particles may be subject to a particle sorting according to particle size and hydrodynamic behavior (McCave et al., 1995). If distinct particle types with high and low isotope concentrations are decoupled from one another during a dislocation process, this may potentially change the particulate ²³¹Pa α ²³⁰Th and ¹⁰Be β ²³⁰Th ratios in sediments after redeposition, which we can test as part of this study.

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99 **1.3 Approaches by other authors**

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101 102 Generally, data on oceanic 231 Pa and 10 Be are still scarce and beyond that, contrasting results led to discussions about their interpretation (Luo and Ku, 1999; Chase et al., 2002;

103 Chase and Anderson, 2004; Luo and Ku, 2004a, b). Therefore, a further development of 104 105 106 107 108 109 110 111 112 113 114 ²³¹Pa α ²³⁰Th and ¹⁰Be α ²³⁰Th as quantitative proxies needs more investigation. Various studies have investigated the scavenging and fractionation of ^{231}Pa , ^{230}Th , and ^{10}Be with respect to particle phases. These different approaches include field observations of geographical variability (e.g. Walter et al., 1997), scavenging experiments under controlled laboratory conditions (e.g. Geibert and Usbeck, 2004), correlations between radionuclides and sediment phases in surface sediments (e.g. Lao et al., 1992) or sediment cores (e.g. Henken-Mellies et al., 1990), studies on the partitioning between dissolved and particulate form in the water column by collecting suspended or settling particles (e.g. Chase et al., 2002), modeling studies (e.g. Siddall et al., 2005; Dutay et al., 2009; Luo et al., 2010), and finally component specific radionuclide analysis by physical or chemical isolation of particle fractions from sediments (e.g. Bourles et al., 1989; Luo and Ku, 1999; Lal et al., 2006).

115 116 117 118 119 120 121 122 123 The methods for component specific radionuclide analysis on isolated sediment fractions are challenging. The caveat associated with the selective chemical leaching of sediment phases is the risk of biasing due to desorption/adsorption kinetics of reactive nuclides from one phase to another (cf. Robinson et al., 2008). Physical separation of sediment phases by sieving and/or settling bears the problem of incomplete isolation of sediment phases (cf. Luo and Ku, 1999). The method of wet sieving also could produce an artifact by an unfavorable leaching of adsorbed isotopes and their loss to the sieving fluid. As a consequence the component specific radionuclides analyzed on mechanically/chemically isolated components may not be representative of their natural composition.

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125 **1.4 How this study addresses the problems**

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127 128 This paper contributes new data on component specific radionuclides in order to assess their sensitivity to redistribution. For our study, we selected two sediment samples (deglacial, 16

ka) from two neighboring sediment cores ("twin cores") from the Southern Atlantic Ocean, previously described by Kretschmer et al. (2010). While the twin cores are similar in most respects, they differ in the extent of sediment redistribution. The sample selection is also advantageous for our study because these samples are mainly composed of only two phases (opal and lithogenic). Carbonate particles contribute a negligibly small amount (<1 weight- %). We decided to physically separate those two sediment phases by a gentle treatment (i.e. without ultrasonication) with natural seawater so that the artificial chemical and physical influence on radionuclide distribution was reduced to a minimum. Any unfavorable leaching loss of radioisotopes to the sieve/settle fluid was monitored by analyzing the radioisotopes within the fluid. Particle separation was achieved by sieving into four size intervals, followed by a separation through settling velocity into fast and slowly settling particle classes. The fast settling class mainly contains particles of lithogenic source (higher specific density) and the slowly settling class mainly contains particles of biogenic opal (lower specific density). Even though this physical separation was incomplete, the resulting sediment fractions carry radioisotope signatures that are strongly differing. Kretschmer et. al (2010) had studied ²³⁰Th 129 130 131 132 133 134 135 136 137 138 139 140 141 142 143 144 145 146 147 148 149 U and Th in the same size and density fractions of the two cores analyzed in this study, the new results in this manuscript are on 231 Pa and 10 Be. This paper describes and compares the distribution of radioisotopes within the sediment samples and the influence of biogenic opal on the ²³¹Pa^{$/230$}Th and ¹⁰Be^{$/230$}Th ratios. The influence of focusing and winnowing on the isotope ratios and the implication on their use as kinematic proxy for deep water ventilation is discussed.

150 **2. MATERIALS AND METHODS**

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152 **2.1. Sediment core location**

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154 155 156 157 158 159 160 161 162 163 164 165 166 167 168 169 170 171 Two neighboring sediment cores located in the Southern Ocean opal belt, Atlantic sector south of the Antarctic Polar front (APF), equal in water depth, but different in sedimentation rates were selected for comparison. These two gravity cores (PS1768-8 and PS1769-1) were recovered during Polarstern cruise ANT-VIII/3 at 52.5930°S, 4.4760°E, 3299 m water depth and at 52.6125°S, 4.4580°E, 3269 m water depth, respectively. The sediment composition is dominated by diatomaceous mud and ice rafted debris (Gersonde and Hempel, 1990). Despite the only 2.5 km lateral distance between the twin cores, they reveal contrasting accumulation rates. The sediment deposition during the last deglacial at core site PS1768-8 (mass accumulation rate (MAR)=10.1 g cm⁻² ka⁻¹, focusing factor (Ψ)=5.9; Frank et al., 1996) is twice as high as at core site PS1769-1 (MAR=4.4 g cm⁻² ka⁻¹, Ψ =3.2; Kretschmer et al., 2010), which is assumed to result from lateral sediment transport by bottom currents. The age model of core PS1768-8 is constrained by a combination of oxygen isotope stratigraphy, siliceous microfossil biofluctuation stratigraphy and radiocarbon dating (Frank et al., 1996). Age control for core PS1769-1 was obtained by correlating the magnetic susceptibility record to that of the parallel core. Each of the twin cores was sampled at one core depth (PS1768-8 at core depth 166-169 cm; PS1769-1 at core depth 14-16 cm) representing the last deglacial period (16 ka).

172 **2.2. Grain size fractionation**

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174 175 The procedure of the grain size fractionation is reported in detail by Kretschmer et al. (2010). The fractionation of two sediment samples into grain size classes was achieved by wet

176 177 178 179 180 181 182 183 sieving and settling. Natural seawater (sea surface from the North Sea) was used as sieving fluid in order to prevent artifacts such as dissolution of particles or desorption of radionuclides that could result from using deionized water. In order to avoid artifacts from particles contained in the natural seawater, it was filtered through polycarbonate-filters (pore size $1 \mu m$) before using it as a sieving fluid. As this natural seawater contains also radionuclides of Th, Pa, and Be, which may produce artifacts to the results, the concentration of radionuclides in the sieving fluid was monitored before and after sediment sieving, as described in section 2.2.2.

184 185 186 187 188 189 190 191 192 193 194 195 196 197 198 199 200 201 The two sediment samples were suspended in filtered seawater and agitated for 2 minutes. The samples were separated into particle size-classes by wet sieving using three nylon sieves (mesh 125 μm, 63 μm, 20 μm). Each of the resulting particle size-classes consisted of particles with different specific densities, i.e. biogenic silica particles (low density) and lithogenic particles (high density). The size and specific density of a particle determine its settling velocity and hence its hydrodynamic behavior in a natural marine system. Therefore, the low density biogenic opal particles were separated from the high density lithogenic particles by their settling velocity in seawater, so that the three sieve fractions 20-63 μ m, 63-125 μ m, and >125 μ m were each split into a "slowly settling" and a "fast settling" fraction. Two settling velocities were specified: the maximum terminal velocity (fast sinking particles) was estimated assuming the average density of quartz $(2.65g \text{ cm}^3)$, the minimum terminal velocity (slowly sinking particles) was determined empirically by measuring settling times of the slowest particles in a settling tube with seawater. Shortly after the fast particles settled down to the bottom of the settling tube the slowly particles were siphoned off. The fraction that rapidly settled to the bottom was again suspended and the separation process was repeated (15-20 times) until the supernatant fluid appeared clear. The separation of the $\langle 20 \mu m\text{-}$ fraction into clay-sized particles ($\langle 2 \mu m \rangle$) and silt-sized particles $(2-20 \mu m)$ was impossible when using seawater, as the smallest particles were

flocculating due to the presence of cations in seawater. For further fractionation into the classes $\langle 2 \mu m \rangle$ and 2-20 μm based on the settling velocity principle the $\langle 20 \mu m \rangle$ -fraction was resuspended in purified water. After a certain settling time (calculated by Stokes' law), the clay-sized fraction $\left($ <2 μ m) remaining in suspension was siphoned off, while the fine silt fraction $(2-20 \mu m)$ settled out from suspension to the bottom of the settling tube and was again resuspended in purified water. This procedure was repeated 10-20 times until the supernatant fluid achieved a low turbidity. Due to the repeated decantations the clay-sized fraction $\langle 2 \mu m \rangle$ was recovered in a large volume of water. As flocculation reagent, a CaCl₂-202 203 204 205 206 207 208 209 210 211 212 solution was added for extracting all particles from the water. The flocculated particles were allowed to settle down for four days, and then supernatant was decanted and stored separately in a canister for further processing (section 2.2.2).

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214 *2.2.1 Characterization of the grain size fractions*

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216 217 218 Further analyses of the sediment fractions include the measurement of biogenic opal (bSi), specific surface area (SSA), and element concentrations. Results were reported by Kretschmer et al. (2010) and are listed in the appendix in Tables A2 and A3.

219 220 221 222 223 224 225 226 The determination of biogenic opal (bSi) followed the description of Müller and Schneider (1993) where opal was extracted with NaOH (1 M) at 85^oC, and dissolved silica was continuously analyzed by molybdate-blue spectrophotometry. The main characteristic of the slowly settling particles is its high concentration of biogenic opal (mainly diatom tests, 75-82 wt-%). Also the fine silt (2-20 μ m) and the clay-sized (<2 μ m) particle fractions contain high amounts of opal (mainly fragments of diatom tests, 76-80 wt-% respectively 61- 68 wt-%). In contrast, the fast settling particles with lower opal content (3-17 wt-%) contain mainly ice rafted rock fragments, mixed with radiolarians and few diatoms.

250 was used for the wet sieving process (i.e. "unused" seawater) and after the wet sieving of the

251 sediment samples. The seawater used for wet sieving was collected in canisters (between 11 L

252 and 22 L per sieved sample, Table 1). Prior to analysis, all seawater samples were acidified

(pH 2-3) by addition of nitric acid (HNO₃). The yield tracers ²²⁹Th, ²³³Pa, ²³⁶U, ⁹Be, and a solution of iron chloride were added. After time for equilibration (24 h) ammonium hydroxide (NH 4-OH) was added to adjust a pH 8-9 for co-precipitating the radionuclides with iron 253 254 255 256 257 258 259 260 261 262 263 264 265 266 267 268 269 270 271 272 273 274 hydroxide (Fe(OH)₃). The precipitate was re-dissolved in $HNO₃$ and ion exchange chemistry was performed following the same protocol as for the sediment samples (section 2.3). The results in Table 1 show that after sediment sieving the radionuclide concentrations in the sieving fluid were higher than before sieving. The difference is assumed to be the amount that was leached from sediment during the sieve/settle process. It accounts for 0.4-2.4 % relative to the total amount of 230 Th, 231 Pa and 10 Be within the respective bulk sample. 2.2.3 *Recovery of* ²³⁰Th, ²³¹Pa, and ¹⁰Be during sediment sieving The isotope recoveries after sediment sieving are calculated by summing up the 230 Th, 231 Pa, and 10 Be concentrations of all particle size fractions and dividing by the concentration of the respective bulk sample. Recoveries of ^{230}Th , ^{231}Pa , and ^{10}Be in sediment fractions sum up to, respectively, 87-90 %, 109-110 %, and 103-104 % relative to the bulk sample (Table 1). **2.3. Radioisotope analyses** The analyses of 231 Pa, 230 Th, and 10 Be were performed in the laboratories of the Marine

275 Geochemistry department at Alfred-Wegener-Institute, Bremerhaven, Germany and the

276 NERC Cosmogenic Isotope Analysis Facility (CIAF) and AMS laboratory, both at the

277 Scottish Universities Environmental Research Center (SUERC), East Kilbride, UK. These

278 laboratories participated in the GEOTRACES inter-calibration program for ^{231}Pa , ^{230}Th and

¹⁰Be (Henderson et al., 2007). The data on ²³¹Pa (T_{1/2}=32.5 ka), ²³⁰Th (T_{1/2}=75.4 ka), and ¹⁰Be 280 $(T_{1/2}=1.5 \text{ Ma})$ reported in this manuscript are decay-corrected for the time of deposition. All 279 281 data are available at [doi:10.1594/PANGAEA.759976.](http://doi.pangaea.de/10.1594/PANGAEA.759976)

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- 283 *2.3.1. 231Pa sample preparation and ICP-MS*
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285 286 287 288 289 290 291 292 293 294 For the isotope dilution analysis by Inductively Coupled Plasma-Sector Field-Mass Spectrometry (ICP-SF-MS) the samples were spiked with ²²⁹Th, ²³³Pa and ²³⁶U. The protactinium spike 233 Pa was produced from a 237 Np solution by separating 233 Pa from its progenitor by chromatography on a silica gel column. The spikes 229 Th, 233 Pa and 236 U were calibrated against the reference standard material UREM 11, an uranium ore. Hansen and Ring (1983) established that this material is in state of radioactive equilibrium. The consensus value for the uranium concentration in UREM 11 is 58.9±0.5 ppm (Hansen and Ring, 1983). From this value, we deduced the concentrations of ²³⁸U (58.48±0.50 ppm), ²³⁴U (3.16±0.03 ng g^{-1}) and ²³⁵U (418.8 \pm 3.6 ng g^{-1}) and calculated the concentrations of the respective daughter nuclides ²³⁰Th (957.3 \pm 8.2 pg g⁻¹) and ²³¹Pa (19.16 \pm 0.16 pg g⁻¹).

295 296 297 298 299 300 301 302 303 304 Samples (10-100 mg) and spikes (0.7 pg ²³³Pa, 9 pg ²²⁹Th, and 800 pg ²³⁶U) were weighed into Teflon vials and fully dissolved in HCl, $HNO₃$, and HF by microwave-assisted digestion (CEM Mars Xpress). Samples were evaporated to dryness, re-dissolved, and coprecipitated with $Fe(OH)₃$. Separation of Pa, Th and U was achieved by ion exchange chromatography with two different exchange resins (Eichrom UTEVA® and BioRad AG 1- $X8@$). The UTEVA columns were conditioned with three column volumes (cv) $HNO₃ (3M)$, then loaded with the samples and rinsed with $HNO₃$ (3cv 3M). Th was eluted with HCl (1cv 9M and 2cv 5M suprapur®) followed by elution of Pa and U with 3cv HCl (0.02M) + HF (0.02M suprapur®). The eluates were collected and evaporated in Teflon beakers. The Th fraction was re-dissolved in $HNO₃$ and as the last step diluted to 5 mL $HNO₃$ (1M). The Pa-U

fraction was re-dissolved in HCl (9M). The AG 1-X8 columns were conditioned with HCl (9M) and loaded with the Pa-U fraction. The first rinsing (3cv HCl 9M) cleaned the column from any remaining Th and was discarded. Subsequently it was collected for Pa (3cv HCl 9M + HF 0.14M) and for U (5cv HCl 0.5M). The separated Pa and U fractions were evaporated and re-dissolved in $HNO₃$ twice, and finally diluted to 5 mL $HNO₃ (1M)$. 305 306 307 308 309

310 311 312 313 314 315 316 317 318 319 320 321 Isotopes of Pa, Th and U were analyzed by isotope dilution using ICP-SF-MS (Element2, Thermo Scientific). Samples were injected into the plasma by a desolvation system (Apex Q^{\circledast} , ESI). During the injection of the Pa fraction, Th hydride (²³²ThH⁺) formation may contribute to the mass 233. This was reduced to a minimum by optimizing the sample and Argon gas flow rates and the plasma temperature. In addition, the 232 Th peak tailing contributes to the masses 233 and 231. The external measurement of a ²³²Th standard solution allowed the assessment of these analytical effects. The contribution to masses 231 and 233 by 232 ThH⁺ and 232 Th peak tailing was corrected by measuring the 232 Th in each Pa sample. The instrument mass bias was assessed externally by bracketing samples with a uranium standard solution. Further corrections accounted for 233 U bleeding (traced by 236 U) and procedural blanks. Procedural blanks (including spike contribution) were equivalent to 0.8-3% (231 Pa) and 0.6% (230 Th) of the samples.

322 323 324 325 326 327 328 329 330 The scavenging process is the principal source for 230 Th and 231 Pa in sediments. Additionally, two further sources contribute variable small amounts to the total 230 Th and ²³¹Pa in sediments: ²³⁰Th and ²³¹Pa that is supported by decaying uranium in lithogenic minerals, and by decaying uranium from authigenic source. For particle flux studies, only the scavenged component of ²³⁰Th and ²³¹Pa is of interest, i.e. the total ²³⁰Th and ²³¹Pa measured in sediments must be corrected for the other two components. The scavenged fraction is referred to as "excess" (excess 230 Th or 230 Th_{xs} and excess 231 Pa or 231 Pa_{xs}) and must be decay corrected for the time of deposition which requires an independent chronology for the sediment core (Henderson and Anderson, 2003).

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2.3.2. 10Be sample preparation and AMS

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354 355 356 Samples were converted into sulfate form by adding H₂SO₄ (1 mL 0.5 M, analytical grade) and evaporating to near dryness. H_2O_2 (6 drops 2 wt %) and H_2SO_4 (2 mL 0.04 M) were added and evaporated two times to a syrupy droplet. The sample was re-

377 **3. RESULTS AND DISCUSSIONS**

400 401 402 The fast sinking coarse particles $>20 \mu$ m (opal-poor) are characterized by ¹⁰Be concentrations (0.50 \pm 0.02 - 0.86 \pm 0.03 \times 10⁹ at/g, Figure 1, Table A2 in the Appendix) and ¹⁰Be/²³⁰Th_{xs} ratios (0.21±0.04 - 0.33±0.06 × 10⁹ at/dpm) lower than within all other particle

classes. The slowly sinking coarse particles $>20 \mu$ m (opal-rich) reveal ¹⁰Be concentrations $(0.84 \pm 0.03 - 1.33 \pm 0.04 \times 10^9 \text{ at/g})$ slightly higher than the fast sinking particles, but their 403 404 405 406 407 408 409 ¹⁰Be/²³⁰Th_{xs} ratios (0.49±0.08 - 0.57±0.10 × 10⁹ at/dpm) are the highest of all. In contrast, the fine particles <20 μ m (opal-rich) carry the highest ¹⁰Be concentrations (3.03±0.09 - 3.11±0.10 \times 10⁹ at/g), but their ¹⁰Be/²³⁰Th_{xs} ratios (0.39±0.07 - 045±0.08 \times 10⁹ at/dpm) are equal (within errors) to the slowly sinking classes. ${}^{10}Be/{}^{230}Th_{xs}$ ratios in all particle size classes exceed the production ratio in seawater (0.136-0.170 \times 10⁹ at/dpm).

410 411 412 413 414 415 416 417 418 Within the same core PS1768-8 at similar core depth (156-181 cm) Frank et al. (2000) analyzed the 10 Be concentration by chemical leaching of the sediment and report a 10 Be concentration of $1.23 \pm 0.04 \times 10^9$ at/g. This is half the concentration of what we measured by complete acid digestion (2.5 \pm 0.1 \times 10⁹ at/g, Figure 1). Either a higher ¹⁰Be signal by complete digestion may derive from a lithogenic 10 Be source or the recovery of authigenic 10 Be from marine sediments by acid leaching may be incomplete (Bourles et al., 1989). The cause for the discrepancy of both 10 Be data cannot be clarified here. In the following paragraph however, the 10 Be sources are discriminated by using the 10 Be/Be ratio within the particle fractions.

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420 *3.2.2 Origin of the10Be/Be signal*

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422 423 424 425 426 427 428 It has been shown that dissolved trace metals are trapped within the silica structure of diatoms so that past metal concentrations in surface waters are recorded by sedimentary opal tests (Lal et al., 2006). This idea is supported by the observation of 10 Be depletion in surface waters of the modern Atlantic ACC which is suggested to result from enhanced scavenging by opal particles in the euphotic zone (Frank et al., 2002). Following this line of argument, our slowly settling opal-rich 20-63 µm particle fractions of PS1768 and PS1769 (containing 75 and 82 % biogenic opal, Table A3) largely reflect an authigenic surface water derived 10 Be/Be

signal (respectively 4.26±0.14 and $4.10\pm0.17 \times 10^{-8}$ at/at, Figure 1). Assuming that ¹⁰Be/Be is equal within all opal tests in all size fractions of PS1768 and PS1769 we can calculate the 429 430 431 432 433 434 435 436 437 438 439 440 441 442 443 444 445 446 447 448 449 450 451 452 453 454 ¹⁰Be/Be ratios for pure (i.e. 100 %) biogenic opal to be 5.7 \pm 0.2 and 5.0 \pm 0.2 \times 10⁻⁸ at/at respectively. The fine particle fractions (<20 µm, clay and fine silt) of PS1768 and PS1769 reveal slightly higher ¹⁰Be/Be ratios (5.46±0.19 and 5.54±0.20 \times 10⁻⁸ at/at, respectively) than slowly settling opal-rich >20 um particles, but they contain less biogenic opal (55 % and 58) %, respectively). We, therefore, suspect that a significant portion of 10 Be/Be must have been contributed by the lithogenic component of the fine fraction (i.e. clay and fine silt). Based on the ¹⁰Be/Be ratios calculated for the pure biogenic opal we can calculate the ¹⁰Be/Be ratios in the pure lithogenic fine fraction to be 5.2 ± 0.3 and $6.3\pm0.3\times10^{-8}$ at/at. The deep and bottom water could be a possible source for 10 Be adsorbed onto clay and fine silt. The ratio reported for the modern deep and bottom water of the circumpolar current is about $10-11\times10^{-8}$ at/at (Kusakabe et al., 1987). Regarding the modern oceanography, it has been suggested that only an advection of $10B$ e from the deep Pacific could explain the high $10B$ e concentrations in the deep water of the Atlantic sector of the Southern Ocean (Frank et al., 2002). Furthermore, the sedimentary clay mineral assemblage of the Eastern South Atlantic sediments indicates a long distance transport of lithogenic material originating from Patagonia and the Antarctic Peninsula (Diekmann et al., 1996). Thus we interpret the high 10 Be concentrations found in the <20 um-fraction to be partly derived from the adsorption on clay minerals that are advected by bottom currents from distant sources. However, if 10 Be/Be was, as it is today, even higher in the deep and bottom water than measured in our lithogenic fraction then the lithogenic component may be diluted with a lower¹⁰Be/Be ratio deriving from aeolian dust that may range from 0.1 to 2×10^{-8} , as reported for surface soils (Barg et al., 1997). The absolute 10 Be contribution to the deep water by aeolian dust is insignificant (Wang et al., 1996), acting mainly as a diluting agent. Dust fluxes of 1-5 $g/m^2/y$ have been reported from ODP core 1090 (Martínez-Garcia et al., 2009).

Lal et al. (2006) analyzed 10 Be/Be in cleaned opal samples that were separated from sediments (5.8 ka - 125 ka) from the ODP core 1093 located very close to our core site. The 455 456 457 458 459 460 461 462 463 ¹⁰Be/Be ratios in these opal tests widely range between 0.19 and 50×10^{-8} at/at and the error associated with the 10 Be/Be data can be as high as 20 %. If such large uncertainties are taken into account for our study, the differences of 10 Be/Be between particle fractions as discussed above would be insignificant. Nevertheless, the data of the $\langle 20 \rangle$ um fraction demonstrate that lithogenic fines substantially contribute to the 10 Be/Be signal so that our conclusion is still valid, i.e. that 10 Be at our study site partly derived from sources other than the ocean surface water.

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465 *3.3. Influence on isotope ratios by sediment redistribution*

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467 468 469 470 471 472 473 474 Comparing the rapid with the slow accumulation site (Figure 1, white and black bars, respectively) they mostly reveal identical (within 1σ) isotope concentrations and ratios. Systematic differences between both sites are not observable. Both sites probably receive particles of the same composition that experience the same water masses, owing to their close neighborhood and the similar water depth. We conclude that mass accumulation rates and consequently the total fluxes of Pa, Th and Be may differ widely between two study sites under contrasting focusing conditions on local scale without affecting their respective isotope concentrations and ratios.

475 476 477 478 479 480 To assess the potential size sorting effect on the isotope signal of the sediments Table 2 provides the percentage of isotopes contributed by each size class in relation to the total inventories. The fine grained $\left($ <20 μ m) and slowly settling opal-rich particles are the main contributors to the isotope inventories and therefore they determine the sedimentary isotope ratios. The size fraction $\left($ <20 μ m) contributes 88-91 % of the total amount of the radioisotopes in the sediment samples. The differences in contribution by the $\langle 20 \mu m\text{-}fraction$

are rather small for the three radioisotopes. The opal-rich particle class of size 20-63 µm is also an important carrier for the three radioisotopes. This class holds more $^{231}Pa_{xs}$ (9-10 %) 481 482 483 484 485 than ¹⁰Be (5-7 %) and ²³⁰Th_{xs} (4-6 %). These differences are essential for the change in isotope ratios in the case of particle sorting by bottom currents as shown in the following paragraph.

486 487 488 489 490 491 492 493 494 495 496 497 498 Figure 4 illustrates a simulation of winnowing by progressively removing the fine sediment component and calculating the resulting radioisotope composition. Two scenarios for winnowing are shown in Figure 4. The removal of only the finest fraction $\left(\langle 20 \right| \text{nm} \right)$ dashed line) leads to a slight increase of the ²³¹Pa/²³⁰Th and the ¹⁰Be/²³⁰Th ratios compared to the measured bulk composition. The removal of the finest fraction $\left(\langle 20 \mu \text{m} \rangle \right)$ together with the slowly settling opal-rich particles (solid line) leads to decreasing isotope ratios. Among those two simulations we expect the second one being the more realistic scenario because the finest particles and the opal particles are both quite susceptible to resuspension and transport owing to their hydrodynamic behavior. A substantial removal of the fine and opal-rich particles (reduction of 70 to 80 % relative to the bulk composition) would reduce the ²³¹Pa α ²³⁰Th ratio to a value lower than the production ratio (0.093). ${}^{10}Be/{}^{230}Th$ ratios would also be affected by winnowing, but less pronounced than $^{231}Pa^{230}Th$. Consequently, strong winnowing at our study site could fractionate 230 Th, 231 Pa, and 10 Be.

499 500 501 502 503 504 505 The opposite will happen in case of concentration of the fine and opal-rich particles by a stronger focusing. However, this stronger focusing could not change significantly the isotope ratios because the investigated cores are already strongly determined by the redistributed fines (<20 µm) and slowly sinking opal-rich particles making 87-91% of the sediment (Table 2) so that the sediment setting is much closer to the full-focusing extreme than to the situation of winnowing. That means that a stronger focusing would hardly change the isotope ratios.

506 507 508 509 510 511 512 513 514 515 516 517 518 519 520 521 522 523 524 525 526 527 528 529 However, regarding the sedimentary setting (strong sediment focusing during glacial and interglacial; Frank et al., 1996), events of strong winnowing are rather unusual at our study site. Therefore, an influence of particle sorting on isotope ratios would more probably occur in regions of strong winnowing, such as, e.g., reported for places south of the Subantarctic Front between 46°S and 48°S (Frank et al., 1996). We can conclude that moderate variations in fluxes of particulate Th, Be, and Pa by lateral advection have negligible influence on the sedimentary isotope ratios found at our study site, and that the isotope ratios reflect particle fluxes on local scale independent of the mass accumulation rate. This conclusion is only valid if redistribution occurs syndepositionally and over short distances. In the case of long range particle transport as already indicated by the 10 Be data any reliable location-specific information on past fluxes of certain particle types cannot be provided. *3.4. Influence of opal content on isotope ratios* 231 Pa_{xs}, 10 Be, and 230 Th_{xs} concentrations are positively correlated to the specific surface area, reflecting the adsorptive binding of all three nuclides to the surfaces of particles (Figure 2a). The ratios ²³¹ $Pa_{xs}/^{230}Th_{xs}$ and ¹⁰Be/²³⁰Th_{xs} are positively correlated with the biogenic opal content of the particles (Figure 2b) confirming findings of earlier studies (Taguchi et al., 1989; François et al., 1993; Asmus et al., 1999; Chase et al., 2002; Bradtmiller et al., 2009). In contrast, the isotope ratios show an anti-correlation to Al (Figure 2b). This is expected as Al represents the lithogenic component which is anti-correlated to biogenic opal. We conclude that the capacity of element adsorption is a function of the surface area of the particle, whereas the ability to fractionate between elements depends on the opal content of

530 the particle.

531 The consequence is, that in regions where biogenic opal is dominating the particle flux 532 533 534 535 536 537 538 539 540 $^{231}Pa_{xs}/^{230}Th_{xs}$ cannot reliably indicate neither the ocean ventilation nor the total mass flux. Thus its use as kinematic proxy and paleoproductivity proxy is limited in those situations. This confirms the statements of earlier studies (e.g. Walter et al., 1997; Keigwin and Boyle, 2008; Scholten et al., 2008; Lippold et al., 2009). Such effects of particle composition on sediment $^{231}Pa^{230}Th$ can be taken into account by analyzing the opal content of the sediment (Gherardi et al., 2009; Guihou et al., in press). Luo et al (2010), however, demonstrate by a 2- D model approach, that a change of particle composition in the Southern Ocean could change the Southern Ocean ²³¹Pa sink and consequently the sediment ²³¹Pa $/$ ²³⁰Th in the Atlantic basin.

541 542 543 544 545 546 In order to test whether the fractionation by opal content is stronger for either one of the nuclides ²³¹Pa or ¹⁰Be, we introduce the distribution ratio $D(N/T)$ as a new parameter. As *D*(*N*/Th) is a ratio of percentages, the direct comparison between the Pa-Th- and Be-Thfractionation is possible and independent of units. The percentage (i.e. the fraction:bulk ratio) of the nuclide *N* (²³¹Pa or ¹⁰Be) within a certain particle size class is divided by the percentage of 230 Th within the same particle size class (equation 1).

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$$
D(N/\text{Th}) = [N_{fraction}/N_{bulk}] / [\text{Th}_{fraction}/\text{Th}_{bulk}], \qquad (1)
$$

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550 551 552 553 554 555 556 where *N* is the concentration of either ²³¹Pa_{xs} or ¹⁰Be, and Th is the concentration of ²³⁰Th_{xs}. Deviation of *D*(*N*/Th) from 1.0 reflects adsorption intensities of *N* different from that of Th. *D*(*N*/Th)>1 indicates preferential adsorption and 0<*D*(*N*/Th)<1 indicates less adsorption of *N* relative to Th. The *D*(*N*/Th) ratio is displayed in Figure 3 on a logarithmic scale where bars are sorted from left to right by increasing opal content. As expected, preferential scavenging of Th relative to Pa and Be is found in the opal-poor particles (*D*(*N*/Th)=0.2-0.9), whereas the inverse situation is found within the opal-rich particle classes (*D*(*N*/Th)=1.3-3.3). Within all

particle classes, $D({}^{10}Be/{}^{230}Th)$ is closer to unity than $D({}^{231}Pa/{}^{230}Th)$. That means the fractionation between 230 Th and 10 Be is less sensitive to the opal content than the fractionation 557 558 559 560 561 562 between 230 Th and 231 Pa. This is consistent with results of other investigators suggesting that 10 Be scavenging is sensitive to both opal and lithogenic particles fluxes (e.g. Lao et al., 1992; Frank et al., 2000; Chase et al., 2002), whereas opal is the major phase for Pa scavenging (Chase et al., 2002; Scholten et al., 2005).

563 564 565 566 567 568 569 The size classes <2 μ m and <20 μ m show $D(^{231}Pa^{230}Th)$ values close to unity in spite of their high opal concentrations (55-70 %). These particle fractions contain the highest concentrations of both ²³¹Pa_{xs} (Figure 1) and ²³⁰Th_{xs} (Table A1). Assuming the most extreme case to explain this observation, virtually all ²³¹Pa could be scavenged by opal particles, and virtually all 230 Th could be scavenged by clay minerals. The actual contribution of the individual particle types to total 231 Pa and 230 Th, however, cannot be fully resolved here and must be addressed in further investigations.

570 **4. CONCLUSIONS**

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572 573 574 575 576 577 578 579 580 581 582 Within various studies the biogenic opal was suggested to be the major influencing factor for ²³¹Pa_{xs} and ¹⁰Be scavenging in the water column. Our data show that ²³¹Pa_{xs} and 10 Be remain strongly associated to the opal-rich particles even after burial in the sediment. The fractionation from ²³⁰Th_{xs} by opal-rich particles is more pronounced for ²³¹Pa_{xs} than for 10 Be. In congruence with previous studies we can conclude that opal fluxes in sediments are better recorded by $^{231}Pa_{xs}$ than by ^{10}Be . Our data support the view of Frank et al. (2000) who conclude from observations on the same sediment core (PS1768-8) that the 10 Be flux rate is not a suitable tracer for biogenic particle flux in the Southern Ocean. Owing to its sensitivity to both lithogenic particle flux and biogenic opal flux 10 Be should be used as tracer for total particle fluxes and its application should be restricted to situations where oceanic particle composition does not change significantly over time.

583 584 585 586 The particle specific 10 Be/Be ratios reveal that the opal-rich particles reflect a local ¹⁰Be signal derived from sea surface, whereas the maximum ¹⁰Be concentrations and ¹⁰Be/Be ratios within the fine particles $\langle 20 \mu m \rangle$ point to another ¹⁰Be source, possibly adsorbed to clay, equilibrated with deep water masses, and advected by bottom currents.

587 588 589 590 591 592 Our experiments have shown that biogenic opal fractionates 231 Pa and 230 Th. In regions of high latitudes where biogenic opal dominates the particle fluxes the scavenging efficiency of ²³¹Pa is increased relative to ²³⁰Th. Therefore, ²³¹Pa^{$/230$}Th is not a reliable proxy for total particle flux in the Southern Ocean. However, based on our data we confirm that $^{231}Pa_{xs}$ is as reasonable tracer for opal fluxes into the Southern Ocean sediments as recently applied by Anderson et al. (2009) and Bradtmiller et al. (2009).

593 594 595 Many studies have applied the ²³¹ $Pa_{xs}/^{230}Th_{xs}$ and ¹⁰Be/²³⁰Th_{xs} ratios for paleoceanographic reconstructions on sediments deposited under rapid accumulation, making the assumption that isotope ratios are not influenced by accumulation rate and sediment

transport. Our results partly confirm this assumption, because concentrations and ratios appear to be insensitive to local variations in the focusing intensity. Based on our data, the fractionation of 231 Pa, 10 Be, and 230 Th by transport processes appears unlikely at our Southern 596 597 598 599 600 Ocean study site. A definite conclusion, however, cannot be drawn as we do not have data from a corresponding winnowing site.

601 602 603 604 605 The simulation of winnowing (Figure 4) shows that particle sorting could change isotope ratios. An extreme sediment winnowing (removal of 70-80 % of fines) would remove high ²³¹Pa^{$/230$}Th and ¹⁰Be $/230$ Th ratios that are associated to the hydrodynamically mobile opal-rich sediment fraction whereas the corresponding low isotope ratios are left behind associated to the hydrodynamically stationary sediment fraction.

606 607 608 609 This study shows the potential effect of particle composition and redistribution on sediment $^{231}Pa^{230}Th$ which must be taken into account for the reconstructions of the Atlantic meridional overturning circulation. In case of drift deposits and sediments with high opal concentrations the use of $^{231}Pa/^{230}Th$ as kinematic proxy should be avoided.

610 611 612 613 The compound-specific information about $^{231}Pa_{xs}/^{230}Th_{xs}$ and $^{10}Be/^{230}Th_{xs}$, which we determined here, shows the potential to trace the origin of an isotopic signature better than the previously accessible bulk isotope ratios. This allows to a certain extent to discriminate signals of remote origin and local signals from surface waters.

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939 **Tables**

940

- 941 Table 1: Sieving recoveries and leaching loss of ^{230}Th , ^{231}Pa and ^{10}Be after the size
- 942 fractionation of the sediment samples. The leaching loss is the amount of ^{230}Th , ^{231}Pa and
- 943 ¹⁰Be that was lost by desorption and/or dissolution during sieving/settling and remained in the
- 944 supernatant water, expressed as the percentage of the total initial amount.

Table 2: Percentage contribution by particle size classes. Each size class contributes a certain 946

fraction to the total flux of opal and particulate $^{230}Th_{xs}$, $^{231}Pa_{xs}$, ^{10}Be and Be. Expressed in 947

- 948 percentage, the comparison between parameters is possible and the potential influence of
- 949 particle size sorting on isotope ratios can be assessed.

950

951 fast = particles of high specific density which were isolated by their rapid settling velocity in

952 seawater (mainly ice rafted debris)

953 slow = particles of low specific density which were isolated by their slow settling velocity in

955 **Appendix**

956

957 958 959 960 961 962 963 Table A1: Particle size specific activities of 231 Pa and 230 Th in two sediment samples. The size classes >20 µm are split in two particle types by their settling velocity (fast and slow) in seawater. Data on $xs^{230}Th_0$ are from Kretschmer et al. (2010). The distribution ratio $D(Pa/Th)$ expresses the particle specific distribution of $xs^{231}Pa_0$ relative to the particle specific distribution of $xs^{230}Th_0$ (for explanation see text). Excess $^{231}Pa_0$ and $^{230}Th_0$ activities are decay corrected to the time of deposition at 16 ka. The uncertainty is given as 1σ . AR = activity ratio.

size class	settle velocity	$^{230}\mathrm{Th}$	$xs^{230}Th_0$	$\mathrm{^{231}Pa}$	$xs^{231}Pa_0$	$xsPa_0/Th_0$	D(Pa/Th)				
[µm]			$[dpm g^{-1}]$			[AR]					
sediment core PS1768-8											
bulk		6.45 ± 0.12	6.68 ± 1.13	0.83 ± 0.03	1.10 ± 0.19	0.16 ± 0.04					
\leq 2		11.91 ± 0.20	12.81 ± 2.16	1.20 ± 0.01	1.59 ± 0.27	0.12 ± 0.03	0.75 ± 0.25				
$2 - 20$		1.30 ± 0.03	1.23 ± 0.21	0.49 ± 0.01	0.67 ± 0.11	0.54 ± 0.13	3.31 ± 1.11				
$<$ 20		6.58 ± 0.11	6.94 ± 1.17	1.05 ± 0.02	1.40 ± 0.24	0.20 ± 0.05	1.23 ± 0.41				
$20 - 63$	fast	2.94 ± 0.06	3.12 ± 0.53	0.12 ± 0.01	0.14 ± 0.03	0.05 ± 0.01	0.27 ± 0.10				
63-125	fast	2.50 ± 0.05	2.68 ± 0.45	0.04 ± 0.00	0.03 ± 0.01	$0.01 \; {\pm} 0.00$	0.07 ± 0.03				
>125	fast	1.39 ± 0.03	1.45 ± 0.25	0.06 ± 0.00	0.08 ± 0.01	0.05 ± 0.01	0.34 ± 0.11				
>63	fast	1.75 ± 0.04	1.84 ± 0.31	0.05 ± 0.00	0.06 ± 0.01	0.03 ± 0.01	0.20 ± 0.07				
$20 - 63$	slow	2.20 ± 0.04	2.34 ± 0.39	0.56 ± 0.01	0.77 ± 0.13	0.33 ± 0.08	2.00 ± 0.67				
63-125	slow	2.01 ± 0.04	2.13 ± 0.36	0.53 ± 0.01	0.72 ± 0.12	0.34 ± 0.08	2.05 ± 0.70				
>125	slow	$3.83\; {\pm}0.08$	4.19 ± 0.71	0.35 ± 0.01	$0.46\; {\pm}0.08$	0.11 ± 0.03	0.67 ± 0.23				
>63	slow	2.12 ± 0.04	2.26 ± 0.38	0.52 ± 0.01	0.71 ± 0.12	0.31 ± 0.08	1.91 ± 0.65				
sediment core PS1769-1											
bulk		6.09 ± 0.10	6.67 ± 1.13	0.80 ± 0.02	1.10 ± 0.19	0.16 ± 0.04					
\leq 2		11.52 ± 0.19	12.73 ± 2.15	1.38 ± 0.03	$1.89 + 0.32$	0.15 ± 0.04	0.90 ± 0.31				
$2 - 20$		2.03 ± 0.04	2.14 ± 0.36	0.50 ± 0.01	0.68 ± 0.11	0.32 ± 0.08	1.93 ± 0.65				
$<$ 20		7.02 ± 0.13	7.69 ± 1.30	1.06 ± 0.05	1.45 ± 0.25	0.19 ± 0.05	1.14 ± 0.39				
$20 - 63$	fast	3.91 ± 0.07	4.16 ± 0.70	0.15 ± 0.01	0.19 ± 0.03	0.05 ± 0.01	0.28 ± 0.10				
63-125	fast	2.31 ± 0.06	2.46 ± 0.42	0.05 ± 0.00	0.06 ± 0.01	0.02 ± 0.01	0.15 ± 0.05				
>125	fast	1.90 ± 0.04	2.03 ± 0.34	0.07 ± 0.01	0.09 ± 0.02	0.04 ± 0.01	0.27 ± 0.10				
>63	fast	2.01 ± 0.04	2.15 ± 0.36	0.06 ± 0.01	0.08 ± 0.02	0.04 ± 0.01	0.23 ± 0.08				
$20 - 63$	slow	1.59 ± 0.03	1.74 ± 0.29	0.58 ± 0.02	0.81 ± 0.14	0.46 ± 0.11	2.82 ± 0.96				
63-125	slow	1.81 ± 0.03	2.00 ± 0.34	0.41 ± 0.01	0.57 ± 0.10	0.28 ± 0.07	1.73 ± 0.59				
>125	slow	3.41 ± 0.07	3.79 ± 0.64	0.55 ± 0.04	0.74 ± 0.13	0.20 ± 0.05	1.18 ± 0.41				
>63	slow	2.06 ± 0.04	2.28 ± 0.38	0.43 ± 0.02	0.60 ± 0.10	0.26 ± 0.06	1.60 ± 0.54				

964

965 966 fast = particles of high specific density which were isolated by their rapid settling velocity in seawater (mainly ice rafted debris)

967 slow = particles of low specific density which were isolated by their slow settling velocity in

975 given as 1σ .

976

977 fast = particles of high specific density which were isolated by their rapid settling velocity in

978 seawater (mainly ice rafted debris)

979 slow = particles of low specific density which were isolated by their slow settling velocity in

981	Table A3: Particle size specific ²³² Th, Aluminum (Al), biogenic opal, and specific surface
982	area (SSA) in two sediment samples. The size classes $>20 \mu m$ are split in two particle types
983	by their settling velocity (fast and slow) in seawater. Data are from Kretschmer et al. (2010).

The uncertainty is given as 1σ . 984

size class	settle velocity	232 Th	Al	opal	SSA			
[µm]		$[dpm g^{-1}]$	[%]	[%]	$[m^2 g^{-1}]$			
	sediment core PS1768-8							
bulk		0.522 ± 0.009	2.03 ± 0.03	54.8				
\leq 2		0.842 ± 0.014	2.37 ± 0.03	67.7				
$2 - 20$		0.256 ± 0.004	1.38 ± 0.02	75.6				
$<$ 20		0.538 ± 0.009	2.05 ± 0.03	54.7				
$20 - 63$	fast	0.306 ± 0.005	5.81 ± 0.11	14.8				
63-125	fast	0.244 ± 0.004	7.30 ± 0.18					
>125	fast	$0.181 + 0.004$	6.17 ± 0.08	3.4				
>63	fast	0.201 ± 0.004	6.53 ± 0.12					
$20 - 63$	slow	0.138 ± 0.002	0.57 ± 0.02	74.9				
63-125	slow	0.106 ± 0.002	0.42 ± 0.01	76.0				
>125	slow	0.222 ± 0.005	0.95 ± 0.01					
>63	slow	0.113 ± 0.002	0.45 ± 0.01					
sediment core PS1769-1								
bulk		$0.461 + 0.008$	1.97 ± 0.03	59.8	19.2			
\leq 2		0.756 ± 0.012	2.31 ± 0.03	61.0	56.5			
$2 - 20$		0.239 ± 0.004	1.17 ± 0.01	79.6	13.0			
$<$ 20		0.528 ± 0.007	1.75 ± 0.02	57.7	33.3			
$20 - 63$	fast	0.507 ± 0.008	5.65 ± 0.11	17.5	2.7			
63-125	fast	0.295 ± 0.006	6.96 ± 0.08	5.0	1.6			
>125	fast	0.231 ± 0.004	6.86 ± 0.03	4.2	3.3			
>63	fast	0.248 ± 0.004	6.89 ± 0.05	4.4	2.2			
$20 - 63$	slow	0.091 ± 0.002	0.35 ± 0.00	81.9	11.2			
63-125	slow	0.096 ± 0.002	0.48 ± 0.00	79.1	9.5			
>125	slow	0.210 ± 0.004	0.78 ± 0.00					
>63	slow	0.113 ± 0.002	$0.53\; {\pm}0.00$					

⁹⁸⁵

987 seawater (mainly ice rafted debris)

988 slow = particles of low specific density which were isolated by their slow settling velocity in

⁹⁸⁶ fast = particles of high specific density which were isolated by their rapid settling velocity in

990 **Figure captions**

991

992 993 994 995 996 997 998 999 1000 1001 1002 Figure 1: Concentrations of ²³¹Pa_{xs} and ¹⁰Be and isotope ratios of excess ²³¹Pa/²³⁰Th and $10Be^{230}$ Th in two sediment samples distributed by particle classes [µm]. The bulk sediment was split into size-classes by sieving. Subsequently, particles of different specific densities in size-classes $>$ 20 um were split by their settling velocity in seawater into a rapidly settling fraction (higher density, mainly ice-rafted debris of lithogenic source) and a slowly settling fraction (lower density, mainly biogenic opal of diatoms). The error bars display the analytical uncertainty of 1σ . Grey horizontal lines indicate the production ratio of $^{231}Pa_{xs}/^{230}Th_{xs}$ (0.093) and a range of the production ratio of $^{10}Be/^{230}Th_{xs}$ (0.136-0.170 \times 10⁹ at dpm^{-1}). This range is based on ²³⁰Th production at 3300 m water depth and a global average production of ¹⁰Be at about 16 ka BP between 1.2 and 1.5×10^6 at cm⁻² yr⁻¹ derived from Christl et al. (2007) and Frank et al. (1997).

1003

1004 1005 1006 1007 1008 1009 Figure 2: (a) ²³¹Pa_{xs} and ²³⁰Th_{xs} specific activities [dpm g⁻¹] and ¹⁰Be concentrations [10⁹ at g⁻¹] ¹] in particle classes (error bar = 1 σ) are plotted against the specific surface area (ssa) [m² g⁻¹] of the particles. (b) ²³¹Pa_{xs}/²³⁰Th_{xs} [activity ratio] and ¹⁰Be/²³⁰Th_{xs} [10⁹ at dpm⁻¹] in particle classes are plotted against biogenic opal [%], aluminum concentration [%] and 232 Th specific activity $\lceil \text{dpm } g^{-1} \rceil$ (error bar = 1 σ). Data and methods on specific surface area, biogenic opal content and Al concentrations are reported by Kretschmer et al. (2010).

1010

1011 Figure 3: The distribution ratio *D(N/Th)* (see text for definition) is calculated for each particle

- 1012 fraction. Values of *D* are plotted on a logarithmic scale, where 0<*D*<1 indicates a
- 1013 fractionation that favors ²³⁰Th over the radionuclide *N* (²³¹Pa or ¹⁰Be), and *D*>1 indicates a
- 1014 fractionation that favors the radionuclide $N(^{231}Pa$ or ¹⁰Be) over ²³⁰Th. The samples are

1015 1016 ordered on the x-axis with increasing biogenic opal content from left to right. Numbers above the bars indicate the particle sizes [µm].

1017

1018 Figure 4: Simulating the effect of winnowing on the isotope ratios (a) $^{231}Pa_{xs}/^{230}Th_{xs}$ and (b)

1019 $10\text{Be}^{230}\text{Th}_{\text{xs}}$ for the sediment sample PS1768-8. Starting with the bulk composition, the fine

1020 fraction is progressively removed and the resulting composition is calculated by using grain

1021 size and isotopic concentration data (Table 2). The resulting isotope ratio depends on the type

1022 of particles being removed: removal of only fine particles <20 µm (dashed line) leads to an

1023 increase of isotope ratios, while the combined removal of $\langle 20 \mu m$ -particles and opal-rich

1024 particles (solid line) leads to a decrease of isotope ratios.

Figure 4

