

1 **Frontogenesis and variability in Denmark Strait and its influence**
2 **on overflow water**

3 Michael A. Spall,* Robert S. Pickart, Peigen Lin

4 *Woods Hole Oceanographic Institution, Woods Hole, MA 02543*

5 Wilken-Jon von Appen

6 *Alfred Wegener Institute Helmholtz Centre for Polar and Marine Research, Bremerhaven,*
7 *Germany*

8 Dana Mastropole

9 *Woods Hole Oceanographic Institution, Woods Hole, MA 02543*

10 H. Valdimarsson

11 *Marine and Freshwater Research Institute, Reykjavik, Iceland*

12 Thomas W. N. Haine and Mattia Almansi

13 *Johns Hopkins University, Baltimore, MD 21218*

14 **Corresponding author address: Michael A. Spall, 360 Woods Hole Road, Woods Hole, MA*
15 *02543.*

16 E-mail: mspall@whoi.edu

ABSTRACT

17 A high-resolution numerical model, together with in-situ and satellite ob-
18 servations, are used to explore the nature and dynamics of the dominant high-
19 frequency (one day to one week) variability in Denmark Strait. Mooring mea-
20 surements in the center of the strait reveal that warm water “flooding events”
21 occur, whereby the North Icelandic Irminger Current (NIIC) propagates off-
22 shore and advects subtropical-origin water northward through the deepest part
23 of the sill. Two other types of mesoscale processes in Denmark Strait have
24 been described previously in the literature, known as “boluses” and “pulses”,
25 associated with a raising and lowering of the overflow water interface. Our
26 measurements reveal that flooding events occur in conjunction with especially
27 pronounced pulses. The model indicates that the NIIC hydrographic front is
28 maintained by a balance between frontogenesis by the large scale flow and
29 frontolysis by baroclinic instability. Specifically, the temperature and salinity
30 tendency equations demonstrate that the eddies act to relax the front, while
31 the mean flow acts to sharpen it. Furthermore, the model reveals that the two
32 dense water processes – boluses and pulses (and hence flooding events) – are
33 dynamically related to each other and tied to the meandering of the hydro-
34 graphic front in the strait. Our study thus provides a general framework for
35 interpreting the short timescale variability of Denmark Strait Overflow Water
36 entering the Irminger Sea.

37 **1. Introduction**

38 Transformation of surface waters to dense overflow waters at high latitudes is a fundamental
39 component of the Atlantic Meridional Overturning Circulation (AMOC). Strong air-sea buoyancy
40 forcing in the Nordic Seas converts the warm, subtropical-origin water to cold water that returns
41 equatorward at depth. The newly-ventilated dense water subsequently flows through gaps in the
42 Greenland-Scotland ridge, the largest of these overflows occurring in Denmark Strait (transport at
43 the sill 3.2–3.5 Sv; Harden et al. (2016); Jochumsen et al. (2017)). As the Denmark Strait Overflow
44 Water (DSOW) descends the continental slope into the Irminger Sea its transport nearly doubles
45 due to entrainment of ambient water, forming the headwaters of the Deep Western Boundary
46 Current (Dickson and Brown 1994). Identifying and diagnosing the dynamical processes that
47 regulate the overflow in Denmark Strait is thus of key importance to improve our understanding
48 of the functioning of the AMOC.

49 It has now been established that there are three pathways of dense water flowing into Denmark
50 Strait: the Shelfbreak East Greenland Current (EGC), the Separated EGC, and the North Icelandic
51 Jet (NIJ, see Fig. 1). The first two currents advect mainly Atlantic-origin overflow water, which
52 is the relatively warm and salty dense water transformed within the rim-current overturning loop
53 of the Nordic Seas (Mauritzen 1996; Våge et al. 2011). The Separated EGC is an offshoot of the
54 shelfbreak EGC that forms near a sharp bend in the bathymetry near 69°N (Våge et al. 2013).
55 By contrast, the NIJ advects predominantly Arctic-origin overflow water that was transformed in
56 the interior of the western Nordic Seas (Våge et al. 2011, 2015). This water is colder, fresher,
57 and denser than the Atlantic-origin overflow water. As the NIJ approaches the strait it merges
58 with the Separated EGC (Harden et al. 2016). The other major current in Denmark Strait is the
59 northward-flowing North Icelandic Irminger Current (NIIC) which advects subtropical-origin wa-

60 ter into the Iceland Sea (Fig. 1). It is believed that the NIIC and NIJ constitute the inflow and
61 outflow, respectively, of a local overturning loop in the Iceland Sea (Våge et al. 2011; Pickart et al.
62 2017).

63 The transport of DSOW (the sum of the branches in Fig. 1) shows no long-term trend and
64 displays little seasonality (Jochumsen et al. 2012, 2017). It has been argued that hydraulic control
65 takes place in the strait, which helps set the mean transport (Whitehead et al. 1974; Whitehead
66 1989; Käse et al. 2003; Nikolopoulos et al. 2003). Whitehead (1989) used shipboard hydrographic
67 data to evaluate the hydraulically-derived volume flux, which gave a transport of 3.9 Sv, in line
68 with the recent mooring estimates noted above. In contrast to the steady nature of the overflow
69 over seasons and years, the flow at the sill is found to vary strongly on short timescales of order 2–
70 5 days (Aagaard and Malmberg 1978; Ross 1978; Macrander et al. 2007; Jochumsen et al. 2017).
71 Earlier studies attributed these high-frequency fluctuations to baroclinic instability (Smith 1976),
72 and fluctuations of a southward-flowing surface current in the strait (Fristedt et al. 1999).

73 Recent analyses of multiple years of hydrographic and velocity data in Denmark Strait have shed
74 further light on the nature of the short timescale variability at the sill. Two different mesoscale
75 processes have been described, which are referred to boluses and pulses. The former (which
76 was first identified decades earlier; Cooper, 1955) corresponds to the passage of a thick lens of
77 cold, dense overflow water. Using a collection of over 100 occupations of the Látrabjarg transect
78 across the Denmark Strait sill over a 23-year period (see Fig. 1 for the location of the transect),
79 Mastropole et al. (2017) found that boluses were present on 41% of the sections. The features are
80 typically found on the western flank of the strait. Using 6 years of mooring data, von Appen et al.
81 (2017) found that boluses are associated with cyclonic azimuthal circulation in the water above the
82 overflow layer as they pass through the strait, although they do not appear to be isolated, coherent

83 eddies. The overflow transport is enhanced when a bolus goes by, mainly due to the raising of the
84 interface between the dense water and the ambient fluid above.

85 The second dominant mesoscale feature found in Denmark Strait is referred to as a pulse (these
86 have only recently been identified, von Appen et al. (2017)). In contrast to boluses, pulses corre-
87 spond to a thinning and acceleration of the overflow layer. Furthermore, von Appen et al. (2017)
88 determined that they are associated with an anti-cyclonic azimuthal flow in the water above the
89 overflow layer. Using the same set of shipboard sections employed by Mastropole et al. (2017), it
90 was demonstrated that the passage of a pulse is coincident with a westward migration of the NIIC
91 hydrographic front throughout the water column. von Appen et al. (2017) speculated that there
92 might be a dynamical connection between the deep pulses of dense water and the variability of the
93 NIIC. As with boluses, the transport of overflow water is enhanced when a pulse occurs, except
94 in this case it is due to the large increase in equatorward velocity of the DSOW which more than
95 compensates the thinning of the layer. Based on the combination of shipboard and mooring data,
96 von Appen et al. (2017) concluded that either a bolus or pulse passes through Denmark Strait on
97 average every 2 days, which is of the correct timescale to account for the dominant high-frequency
98 variability noted by the many previous studies.

99 Numerical simulations have also been used to investigate the energetic fluctuations in Denmark
100 Strait. Käse et al. (2003) ran a model with an upstream reservoir of dense water (i.e. a “dam
101 break” problem) to investigate the resulting flow through the strait. Eddies were generated along
102 the path of the dense water which had similar characteristics to boluses. Spall and Price (1998)
103 found that an outflow of dense and intermediate waters through the strait produced strong cyclones
104 that were in general agreement with observations. The main driver of the explosive cyclogenesis
105 was stretching of the intermediate layer above the overflow water. Almansi et al. (2017) analyzed
106 a high-resolution general circulation model with realistic geometry and atmospheric forcing. The

107 variability in Denmark Strait was found to be quite similar to that seen in the observations of
108 Mastropole et al. (2017) and von Appen et al. (2017). In particular, both boluses and pulses
109 were present in the model with similar characteristics and time-space scales. Consistent with the
110 data, the boluses were cyclonic and the pulses were anti-cyclonic, and both features resulted in an
111 increase in equatorward transport of overflow water. Almansi et al. (2017) also determined that
112 sea surface height anomalies were centered upstream of the sill when boluses and pulses crossed
113 the strait, which is consistent with baroclinic instability. It still remains to be determined, however,
114 what the precise dynamics are of both types of features, and if they are related to each other.

115 In the present study we further investigate the nature of the high-frequency variability in Den-
116 mark Strait using a combination of in-situ and satellite data together with a high-resolution nu-
117 merical model (the same model employed by Almansi et al. (2017)). First, we describe another
118 mesoscale process that occurs in Denmark Strait which is referred to as a flooding event. Dur-
119 ing such an event, warm subtropical-origin water flows northward through the deepest part of the
120 sill, associated with a westward migration of the NIIC. Next, using the different data sources, it is
121 shown that the flooding events are in fact related to the previously described pulses. The numerical
122 model is then used to explore the dominant variability in the strait, focusing on the role of the NIIC
123 front; namely, the occurrence of frontogenesis and instability. In doing so, we demonstrate that the
124 three types of DSOW variability – boluses, pulses, and flooding events – are tied together within
125 a single dynamical framework. Our results thus provide insight regarding the time-dependent flux
126 of overflow water into the Irminger Sea.

127 **2. Data and Numerical Model**

128 *a. In-situ Data*

129 The primary in-situ data used in the study are from mooring DS1 deployed in Denmark Strait
130 at the deepest part of the sill (Fig. 1). The mooring contains an RDI 75-KHz upward-facing
131 Long Ranger acoustic Doppler current profiler (ADCP) situated at 648 m, roughly 8 m above the
132 sea floor. Velocity was recorded every hour in 16-m bins over the depth range 80–630 m. The
133 dominant tides were removed using the T_TIDE software package (Pawlocwicz et al. 2002). It has
134 been documented that in some years the DS1 ADCP underestimates the near-bottom velocity due
135 to interference from side-lobe reflections (Jochumsen et al. 2017). However, this does not affect
136 the results presented here as we are not concerned with the detailed structure of the near-bottom
137 flow. In particular, the identification and characterization of the warm water flooding events in
138 Denmark Strait are not qualitatively influenced by this. The same was true for the analysis of
139 boluses and pulses carried out by von Appen et al. (2017). Here we consider the 10-year time
140 period from 2005-15 (although there are no data for the 2006-7 deployment year). Velocities are
141 rotated to the along-stream (positive 230°T towards the Irminger Sea) and cross-stream (positive
142 140°T towards Iceland) directions. The along-stream direction is dictated by the long-term mean
143 flow vector from DS1 (von Appen et al. 2017). We also use the data from the temperature sensor
144 on the ADCP, which has a resolution of 0.1°C and accuracy of 0.4°C . Comparisons with calibrated
145 MicroCATs indicate that the accuracy of the ADCP thermistor is in fact better than this (D. Torres,
146 pers. comm., 2018).

147 *b. Satellite Data*

148 The along-track absolute dynamic topography (ADT) data used in the study were obtained from
149 Copernicus Marine and Environment Monitoring Service (CMEMS, <http://marine.copernicus.eu/>),
150 which provides the satellite altimetry product formerly distributed by Archiving Validation and
151 Interpretation of Satellite Data in the Ocean (AVISO). The measurements were made by the
152 Jason-1 satellite until October 2008, after which Jason-2 became operational. The along-track
153 data have a spatial resolution of 12 km, and, since Denmark Strait is near the latitude of the
154 turning point of the satellite, the temporal resolution is roughly 2 days. We also use the daily
155 gridded surface geostrophic velocity product from CMEMS which merges the multiple satellite
156 altimeter measurements and has a horizontal resolution of 0.25° in longitude and latitude. For
157 sea surface temperature (SST) we employ MODIS Aqua Level 3 imagery with 9 km resolution
158 (<https://podaac.jpl.nasa.gov/>). The time period considered for all datasets is 2006-2015, which
159 roughly corresponds to the time period of the DS1 mooring data used here.

160 *c. Numerical model description*

161 The numerical circulation model dataset used here is a high-resolution realistic run of the Mas-
162 sachusetts Institute of Technology General Circulation Model (MITgcm; Marshall et al. 1997). It is
163 publicly available on the Johns Hopkins University SciServer system (<http://www.sciserver.org/integration/oceanography/>). The model setup is explained in detail in Almansi et al.
164 (2017), but is briefly described here. The model was run for 1 year (from September 2007 to
165 August 2008) assuming hydrostatic balance, implementing a non-linear formulation for the free-
166 surface, and applying the non-local K-Profile Parameterization for vertical mixing. ERA-Interim
167 (Dee et al. 2011) provides the atmospheric boundary conditions used to force the oceanic and sea
168 ice components every 3 hours. The horizontal resolution is 2 km over the region of interest, the
169

170 vertical resolution varies from 1 m at the surface to 15 m below a depth of 120 m, and the numer-
171 ical solutions have been stored every 6 h. This high resolution is appropriate for studies, such as
172 the present one, focusing on high-frequency mesoscale features.

173 *d. Model prior validation*

174 The model hydrography and circulation in Denmark Strait have been previously compared with
175 available observations (Almansi et al. 2017). Overall, the model does an excellent job of capturing
176 the major currents and water masses observed in Denmark Strait. Almansi et al. (2017) identified
177 the subtropical-origin (Irminger) water, the recirculated Irminger water, and both types of overflow
178 water (Arctic-origin and Atlantic-origin) in the model. The currents advecting these water masses
179 to the strait are well captured by the model, and the simulated NIIC and DSOW velocities are
180 similar to the measurements reported by Våge et al. (2011). The properties of the water masses
181 mentioned above are consistent with the historical CTD data analyzed by Mastropole et al. (2017).
182 The model does, however, appear to have a small bias in temperature affecting the density in
183 the deep part of the water column. Specifically, while the isopycnal structure across the strait is
184 very similar to that seen in observations, the measured overflow is slightly denser than the model
185 overflow (the magnitude of density biases does not exceed 0.1 kg m^{-3} , corresponding to a model
186 warm bias of less than 1°C).

187 **3. Warm water flooding events in Denmark Strait**

188 *a. Evolution of a composite event using the mooring data*

189 In line with previous studies, we define DSOW as water denser than 27.8 kg m^{-3} (Dickson
190 and Brown 1994). In the mean, the dense water is banked against the western side of the trough
191 separating the Iceland and Greenland shelfbreaks as it flows through the strait (Fig. 2a). The 27.8

192 isopycnal rises 250 m from east to west, and the coldest, densest Arctic-origin overflow water is
193 found at the bottom of the trough where the DS1 mooring is located. When boluses pass by, the
194 interface can rise to within 200 m of the surface, associated with the thick lens of Arctic-origin
195 water (Mastropole et al. 2017). By contrast, the interface deepens when pulses go by such that the
196 DSOW layer can be less than 100 m thick (von Appen et al. 2017).

197 On occasion, much of the trough at Denmark Strait is filled with warm, sub-tropical origin water.
198 An example is shown in Fig. 2b, where water warmer than 4°C occupies most of the Látrabjarg
199 section, including the deepest part of the sill. During this occupation of the line there was only
200 a small amount of DSOW present. How often does this situation occur? For the collection of
201 111 shipboard sections considered by Mastropole et al. (2017) only a small number of realizations
202 captured this state (the section shown in Fig. 2b is the most pronounced example), suggesting that
203 the condition is not common. To investigate this more definitively we considered 9 years of DS1
204 mooring data.

205 Using a graphical user interface (GUI) applied to the mooring data, we identified all of the
206 instances in which the temperature at the bottom of the trough exceeded 1°C. The majority of them
207 (>70%) were associated with northward flow through the strait. On occasion the temperature was
208 warmer than 6°C. We refer to this condition as a warm water flooding event; i.e., when the bottom
209 temperature is warmer than 1°C and the along-stream flow reverses to the north. Over the 9-year
210 record there were 151 such events; on average, one per month. There is no apparent seasonality or
211 long term trend to the flooding events (Fig. 3).

212 The GUI revealed that there was a well-defined, consistent evolution associated with this pro-
213 cess. To quantify this we constructed a composite average of all the events, aligning each one at
214 the time of maximum temperature. Fig. 4 shows the resulting composite timeseries of temperature
215 and velocity over a ± 3 -day period surrounding the peak temperature, which is defined as time = 0.

216 The top panel is bottom temperature, the next two panels are depth-dependent along-stream and
217 cross-stream velocity. The fourth panel is the flow averaged over the approximate depth range of
218 the DSOW layer, and that averaged in the water column above this. (The final panel is discussed
219 below.)

220 Prior to the onset of the flooding event, the DSOW is flowing to the southwest as it does in the
221 mean (see Fig. 1), and the bottom temperature is colder than 0°C corresponding to the Arctic-origin
222 overflow water. Roughly a day before the peak of the event, the flow reverses to the northwest,
223 reaching maximum strength ~ 12 hours before the temperature attains its highest value (3.4°C in
224 the composite mean). Note that the northwest flow is surface-intensified, which suggests that it is
225 the NIIC. As the temperature falls, the southwest flow of DSOW is established again. On average
226 the events last 1.2 days, with the velocity signal leading the temperature signal.

227 *b. Sea surface signature of an event*

228 To shed further light on the nature of the flooding events, we analyzed the along-track ADT of
229 the sea surface in the vicinity of Denmark Strait from CMEMS (see section 2b). There are four
230 satellite tracks that cross the strait, passing almost directly over the DS1 mooring site (Fig. 5).
231 We considered the portion of the tracks within the rectangle in Fig. 5, and computed the surface
232 geostrophic velocity associated with each crossing (using the component of velocity in the cross-
233 track direction). Following this, we constructed composites of the surface velocity for each day of
234 the flooding events, covering a ± 3 day period centered on the peak of the event. The times of the
235 events were identified from the mooring data.

236 This calculation reveals that the flooding events are indeed associated with a westward propa-
237 gation of the NIIC (Fig. 6a). Prior to the event the NIIC is located over the outer Iceland slope
238 and the velocity is equatorward in the trough. As the event progresses, the NIIC moves across the

239 strait at a rate of $\sim 20 \text{ km d}^{-1}$ (indicated by the dashed line in Fig. 6a). At the end of the event
240 the surface velocity at the DS1 site becomes equatorward again, consistent with the mooring data.
241 The northward-flowing NIIC is also re-established over the Iceland slope.

242 This sequence is confirmed using independent SST satellite data from MODIS. We constructed
243 the analogous composite of the SST gradient across the strait near the mooring location (Fig. 6b).
244 Although the SST signature is more noisy than the ADT signal, the NIIC is clearly identifiable as
245 a maximum in SST gradient, i.e. the hydrographic front of the warm, salty Irminger water. One
246 sees that the front propagates westward in conjunction with the surface velocity signal during the
247 flooding event at the same rate of 20 km d^{-1} . The extent of the frontal excursion is consistent with
248 that deduced from the time integral of the depth-mean cross-stream velocity of the upper layer
249 using the mooring data (roughly 40 km, see Fig. 4e).¹

250 *c. Relationship of flooding events to pulses*

251 As described by von Appen et al. (2017), the pulses in Denmark Strait are associated with an
252 increased equatorward flow of DSOW in concert with a change in cross-stream velocity from
253 negative to positive; i.e., prior to the pulse the cross-stream flow is towards Greenland, and subse-
254 quent to the pulse it is towards Iceland, resulting in an anti-cyclonic sense of rotation. Throughout
255 the event the cross-stream velocity signal is surface intensified. At the same time, the NIIC hydro-
256 graphic front is displaced westward during a pulse. This led von Appen et al. (2017) to hypothesize
257 that the occurrence of pulses was dynamically related to the behavior of the NIIC.

258 In the composite flooding event of Fig. 4, the along-stream velocity in the overflow experiences a
259 temporary peak roughly one day after the temperature peak. In addition, the cross-stream velocity
260 in the water column above the overflow changes sign such that, prior to the along-stream peak,

¹It is unclear why the surface signature of the NIIC does not appear to propagate back to the east after the flooding event.

261 the upper-layer flow is towards Greenland, while afterwards it is towards Iceland. Furthermore,
262 the overflow interface height decreases during this sequence of events.² The combination of these
263 signals is strikingly reminiscent of the pulses described by von Appen et al. (2017) (see their Fig.
264 9). Our analysis thus implies that flooding events are in fact related to pulses and do not represent
265 a different type of mesoscale process in Denmark Strait.

266 We note that not every flooding event detected by the DS1 mooring was followed by a pulse.
267 One possible explanation for this is that, for especially strong flooding events (associated with
268 pronounced excursions of the NIIC), the subsequent pulse occurred to the west of the mooring
269 location. This is consistent with the fact that these events were associated with warmer bottom
270 temperatures. Nonetheless, the majority of the flooding events were followed by a pulse (as is
271 evident from the composite of Fig. 4). However, the opposite is not true. The mooring data
272 indicate that not all pulses are preceded by warm water flooding the deepest part of the trough. In
273 fact, there are many more pulses than there are flooding events. von Appen et al. (2017) determined
274 that pulses occur on average every 5.4 days, whereas flooding events take place roughly once a
275 month (Fig. 3b). The likely explanation for this is that flooding events measured by the mooring
276 are simply cases when the NIIC is displaced far enough westward to reach the mooring site in the
277 trough. This is supported by the numerical model results presented below.

278 **4. Frontogenesis**

279 We now focus on the structure of the boundary current system in the vicinity of Denmark Strait
280 using the model data. Consider first a meridional section at 25.75°W, north of the Látrabjarg line
281 (the section is shown in Fig. 7d). The three major currents near the Iceland slope are evident: the

²When the bottom temperature is above 0°C there is no overflow water present, and, consequently, there is no interface between the overflow layer and the water above. This explains the gap in the interface time series.

282 Separated EGC, the NIJ, and the NIIC (Fig. 7a, the former two are in the process of merging at this
283 point). The mid-depth maximum in the southwestward flowing NIJ is supported by the change in
284 slope of the isopycnals around 500 m depth. The northeastward flowing NIIC is strongest near the
285 bottom, with the vertical shear balanced by the upward sloping isopycnals near the surface. The
286 temperature and salinity sections delineate the different water masses transported by these three
287 velocity cores (Figs. 7b, c). The water in the NIIC is warm and salty, while the overflow water in
288 the NIJ is colder and slightly fresher. The near-surface portion of the Separated EGC transports
289 a combination of cold and very fresh polar water alongside Atlantic water. In the upper part of
290 the water column the density is controlled by salinity, while near the bottom it is controlled by
291 temperature.

292 *a. Energetics*

293 The eddy kinetic energy at 50 m depth and at 420 m depth are shown in Fig. 8. Eddies here are
294 defined as deviations from the time mean fields so include variability at all frequencies less than 1
295 year. At both levels local maxima of $O(0.1 \text{ m}^2 \text{ s}^{-2})$ are found near and south of the Denmark Strait
296 sill. This is larger than the estimate by Håvik et al. (2017) of $0.02 \text{ m}^2 \text{ s}^{-2}$, which was based on
297 along-track sea surface height satellite measurements and is thus probably an underestimate due
298 to limited spatial and temporal resolution of the data. The upper layer shows a band of enhanced
299 variability extending to the northeast while the deeper level is enhanced to the southwest. There
300 is also enhanced eddy kinetic energy over the shelf west of Denmark Strait, which is a distinct
301 feature and will not be discussed further. The dominant variability lies along the 600 m isobath
302 both to the northeast and southwest of Denmark Strait.

303 The correlation between the high frequency temperature variability at 50 m depth and that at
304 420 m depth is shown in Fig. 9. The time series of temperature at each grid point has been high

305 pass filtered for periods less than 10 days to highlight the vertical coherence of the mesoscale
 306 variability and filter out seasonal and lower frequency influences. The correlation was computed
 307 at zero time lag; positive and negative time lags between the shallow and deep time series produce
 308 lower correlations. There is a band of high correlation (exceeding 0.8) extending from Denmark
 309 Strait to the southwest. This is the signal of the strong cyclones that are known to be generated due
 310 to the descending overflow of dense water south of the strait (Smith 1976; Spall and Price 1998;
 311 Käse et al. 2003). This is also reflected in the eddy kinetic energy fields. However, the correlation
 312 northeast of the strait is much lower even though the near-surface kinetic energy there is similar to
 313 that found southwest of the strait.

314 The energy source of the variability in the vicinity of Denmark Strait is now diagnosed. Sources
 315 due to internal instabilities are characterized as either baroclinic or barotropic, depending on
 316 whether the eddy energy is derived from the mean potential energy (BC=baroclinic) or the mean
 317 kinetic energy (BT=barotropic). The energy conversion rates are calculated as

$$BC = -g/\rho_0(\gamma_x \overline{u'\sigma'_\theta} + \gamma_y \overline{v'\sigma'_\theta}) \quad BT = -\overline{u'v'}(\overline{u}_y + \overline{v}_x) \quad (1)$$

318 where u and v are the zonal and meridional velocities, overbars indicate the time average, σ_θ is the
 319 potential density, primes are deviations from the time mean, g is gravitational acceleration, and ρ_0
 320 is a reference density. Mean potential energy is converted to eddy energy by the horizontal eddy
 321 density flux in the direction of the mean isopycnal slope, where γ_x and γ_y are the isopycnal slopes
 322 in the zonal and meridional directions (BC). Mean kinetic energy is converted to eddy energy by
 323 the eddy momentum flux $\overline{u'v'}$ across the mean horizontal velocity shear $\overline{u}_y + \overline{v}_x$ (BT). Positive
 324 values indicate a transfer of energy from the mean fields to the eddy fields.

325 The baroclinic conversion rate BC at 420 m depth and 50 m depth are shown in Fig. 10a and b.
 326 At depth there is energy extraction from the mean in a band extending from the sill towards the

327 southwest, along the path of the overflow water. At 50 m depth there is also a positive conversion,
328 although it is more spatially variable and largest from the sill towards the northeast, along the
329 path of the NIIC. Both baroclinic instability and symmetric instability result in $BC > 0$, but the
330 conditions for symmetric instability – that Ertel potential vorticity be negative – are not satisfied
331 outside of $O(10m)$ thick surface and bottom boundary layers. Hence we attribute the source of the
332 variability to baroclinic instability of both the dense overflow waters and the NIIC. The NIJ does
333 not exhibit significant energy conversion upstream of the sill even though there is baroclinic shear
334 present. The energy conversion terms and vertical coherence suggest that two distinct forms of
335 variability are present: a coupled mode in and south of the strait, and a surface intensified mode in
336 the NIIC along the northwest Iceland shelfbreak.

337 The vertical structure of BC and along-strait velocity at the sill are shown in Figs. 10c, d for
338 the model equivalent of the Látrabjarg transect (red line in Fig. 10a, b). There are two cores of
339 energy conversion, one near the bottom and one near the surface. The bottom region is extracting
340 energy from the sloped isopycnals associated with the weakly stratified, dense overflow water. The
341 upper region is extracting energy from the density gradient between the lighter southward-flowing
342 water in the Separated EGC and denser northward-flowing water in the NIIC. The isopycnals are
343 relatively flat in the middle of the water column and so provide little source for energy extraction
344 there. The barotropic conversion term BT was also calculated but was found to be generally
345 much smaller, especially when integrated across the current because there are regions of offsetting
346 positive and negative eddy momentum fluxes. The Reynolds stresses act primarily to shift the
347 location of the front slightly.

348 *b. Temperature and salinity balances*

349 The variability associated with the baroclinic conversion term acts to relax the isopycnal slopes.
350 In terms of the temperature tendency equation, this causes the warm side of the front to cool and
351 the cold side of the front to warm. The contribution to the tendency of temperature and salinity
352 by the mean and the negative of the eddy advection terms³ at 50 m depth are shown in Fig. 11.
353 The mean advection is making the northwest side of the front colder and fresher and the southeast
354 side warmer and saltier. The eddy advection terms generally counteract the mean flow, making the
355 northwest side of the front warmer and saltier and the southeast side colder and fresher. Although
356 there is some spatial variability, this general balance is found all along the frontal region.

357 A similar balance is found at all depths north of Denmark Strait. This is seen in the vertical
358 sections of temperature and salinity tendency for the transect along 27.75°W (Fig. 12). Although
359 the balance is not exact, to leading order the flow is adiabatic with the mean flow acting to increase
360 the lateral temperature and salinity gradients and the eddies acting to weaken the front. Note that
361 the isopycnal slope changes sign with depth. Therefore, in the upper ocean the mean advection
362 of salinity is acting to increase the horizontal density gradient, while at depth the advection of
363 temperature is acting to increase the horizontal density gradient.

364 The mean model SST and surface velocity field in the region encompassing Denmark Strait
365 are shown in Fig. 13a. This reveals the convergent mean flow that supports the hydrographic
366 front: cold, fresh surface waters are advected from the north, and warm, salty surface waters are
367 advected from the south. The analogous fields derived from the satellite ADT and SST data are
368 shown in Fig. 13b. While the model produces a sharper temperature front, likely due to the coarser
369 resolution and processing of the satellite data, the overall surface velocity and temperature fields
370 are very similar between the model and observations. It is noteworthy that the mean velocity field

³The negative is shown in order to make it easier to compare with the mean.

371 at the surface is broadly distributed across the strait; it is not confined to narrow boundary currents.
372 Its significant orthogonal relationship with surface density (and temperature and salinity) gradients
373 points to the importance of the barotropic component of the flow.

374 A measure of the influence of the mean flow on the tracer fields is given by the strain field ε ,
375 defined as

$$\varepsilon = [(u_x - v_y)^2 + (v_x + u_y)^2]^{1/2}, \quad (2)$$

376 where subscripts indicate partial differentiation and u, v are the mean velocities. As seen in
377 Fig. 13c, the time-mean model strain is largest along the boundary between waters emanating
378 from the north and those originating from the south.⁴ It is clear that the mean velocity acts to in-
379 crease the horizontal gradient of temperature along this boundary (also the gradients of salinity and
380 density). The strain calculated from the satellite-derived velocity field (Fig. 13d) shows a similar
381 pattern with a maximum along the temperature front northeast of Denmark Strait (the magnitude
382 is smaller, which is expected given the low resolution of the gridded velocity field). Interestingly,
383 calculation of the strain over the Faroe-Bank overflow region does not show a similar enhance-
384 ment; this appears to be unique to the Denmark Strait. This is likely because Denmark Strait lies
385 on the western boundary of the Nordic Seas and is thus the location of southward flowing, low
386 salinity waters, and the topographic configuration steers the northward and southward flowing wa-
387 ters through the narrow strait. These waters were identified as a key driver of the cyclogenesis
388 south of Denmark Strait by Spall and Price (1998).

389 Eddy fluxes play a leading role in the temperature, salinity, and density budgets by acting to
390 adiabatically balance the mean flow. This is important because, if the eddy field is not suffi-
391 ciently resolved, the mean flow will collapse the gradients to such small scales that parameter-

⁴There is also very large strain over the Djupall Canyon near 66.5°N, 24°W, perhaps also weakly present in the satellite data.

392 ized or numerical lateral mixing will become important (Spall 1997; McWilliams and Molemaker
393 2011). Such subgridscale mixing often artificially introduces diapycnal mixing, modifying the
394 water masses and the resulting transports of heat and salt.

395 *c. Relation to overflow water variability*

396 Returning to the short timescale variations of the DSOW, we are now in a better position to
397 understand the underlying cause of these fluctuations and relate the boluses and pulses to each
398 other (it has already been demonstrated that flooding events are extreme versions of pulses). The
399 model has revealed that the NIIC is baroclinically unstable. A manifestation of this is the mean-
400 dering of the hydrographic front (akin to the Gulf Stream north wall). When flooding events are
401 present, meanders of the NIIC projected onto the nearly zonal satellite tracks in Fig. 5 produce a
402 zonal phase speed of approximately 20 km d^{-1} to the west, very close to the satellite observations
403 (see Fig. 6). The meandering results in enhanced eddy kinetic energy at periods of several days,
404 consistent with the observed overflow variability measured at the sill.

405 Using the model fields we composited the SST during periods when flooding events were present
406 at the sill and periods when boluses of DSOW were present at the sill. Flooding events were de-
407 fined as times when the temperature at 600 m depth was warmer than 3° C and boluses were de-
408 fined as times when the temperature at 400 m depth was colder than 1° C . While these definitions
409 differ slightly from those used above for the mooring, the model provides more information in the
410 vertical and these choices allow for a clearer identification of these events. As seen in Fig. 14,
411 flooding events occur during meander crests (northwestward excursions of the NIIC front), while
412 boluses occur during meander troughs (southeastward excursions of the NIIC front). This is con-
413 sistent with respect to the observed lateral movement of the front during pulses (von Appen et al.
414 2017) and during flooding events (Fig. 6). (There is no observational evidence to date of shore-

415 ward excursions of the NIIC during boluses, mainly because of a seasonal bias in the collection of
416 Látrabjarg transects, see von Appen et al. (2017). However, in the model of Almansi et al. (2017)
417 the NIIC moved onshore during bolus events.) The observed cyclonic versus anti-cyclonic sense
418 of rotation of the boluses versus pulses is also consistent with the meander troughs versus crests
419 seen in the model composites. These results thus link the two dominant modes of observed over-
420 flow variability to a single dynamical process associated with the instability of the hydrographic
421 front in Denmark Strait.

422 **5. Discussion and Summary**

423 In-situ observations, remotely sensed data, and a regional high resolution numerical model have
424 been used to provide a unifying view of high frequency variability in the vicinity of Denmark
425 Strait. The observed flooding of warm, salty northward-flowing water through the deepest part of
426 the strait is shown to be associated with a westward shift of the NIIC. These flooding events occur
427 about once per month and appear to be extreme versions of the more common and previously
428 described anticyclonic pulses of dense water.

429 It was also shown that the front separating the northward-flowing NIIC from the southward-
430 flowing Separated EGC / NIJ (which are essentially merged in the strait) is baroclinically unstable.
431 There are two dominant regions of energy conversion which act to flatten the isopycnals: one in
432 the upper layer and one near the bottom. The large-scale mean flow in both the numerical model
433 and that inferred from sea surface height data are broadly southward-flowing north of the sill and
434 northward-flowing south of the sill. The water north of the sill is cold and fresh while the water
435 south of the sill is warm and salty. This provides a confluent flow that acts to sharpen the horizontal
436 gradients of temperature and salinity, and to steepen the isopycnals throughout the water column.
437 Over the long-term mean, the tendency of the mean flow to steepen the front is nearly adiabatically

438 balanced by the tendency for eddies generated by baroclinic instability to relax the front. Thus we
439 view the high frequency (one day to one week) variability in the vicinity of Denmark Strait to
440 derive from a baroclinic front maintained locally by the large-scale mean flow. Similar balances
441 between confluent flows and frontal instability have been discussed in the context of submesoscale
442 upper ocean fronts (Spall 1997; McWilliams and Molemaker 2011).

443 Our results suggest that the intense growth of cyclones southwest of the sill results at least in
444 part from the localization of the confluent flow in the vicinity of the sill. Once the strong baroclinic
445 shear that is formed near the sill is free from the frontogenetic effect of the large-scale mean flow,
446 baroclinic instability can grow unchecked. The growth rate is likely also enhanced by the large
447 horizontal gradients in potential vorticity found south of the sill (Spall and Price 1998). A similar,
448 but weaker, region of baroclinic conversion is found in the upper ocean to the northeast of the
449 sill. There does not appear to be an analogous confluence flow in the vicinity of the Faroe Bank
450 overflow or the Mediterranean overflow, but there is a similar confluence with strong fronts and
451 enhanced eddy variability in Fram Strait (Hattermann et al. 2016).

452 This highlights the importance of properly representing both baroclinic instability and eddy
453 fluxes on small scales. If a model is unable to represent the energy conversion and growth of
454 eddies and meanders, it will not be able to properly arrest the frontogenetic effect of the con-
455 fluent flow. Eventually the front will sharpen to the point where numerical mixing balances the
456 mean flow. This will introduce diapycnal mixing that dilutes the water masses, both those flowing
457 northward into the Nordic Seas and those flowing southward forming the headwaters of the Deep
458 Western Boundary Current. In addition to producing water masses of incorrect density, this mixing
459 will alter the heat and freshwater transports associated with the AMOC. This is analogous to the
460 “Veronis Effect” (Veronis 1975), previously identified as a major source of error in the meridional

461 heat transport in climate models that results from numerical diapycnal mixing in the vicinity of
462 the Gulf Stream.

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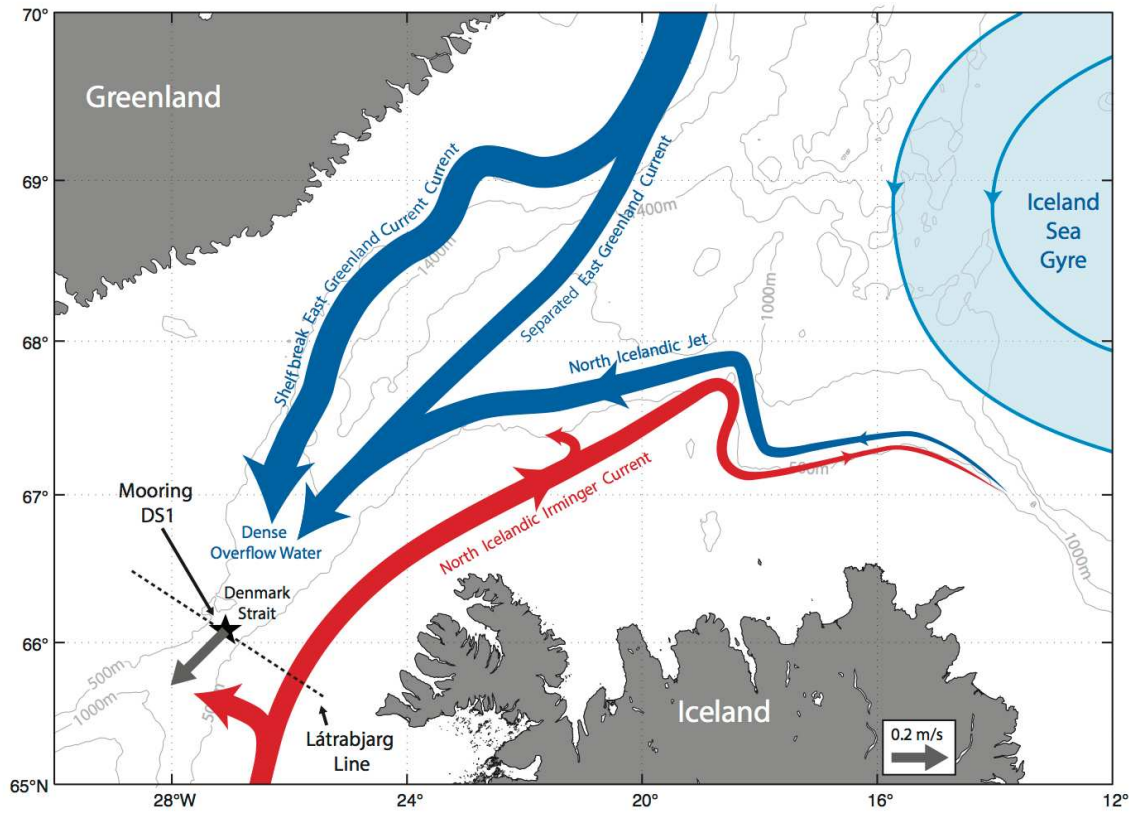
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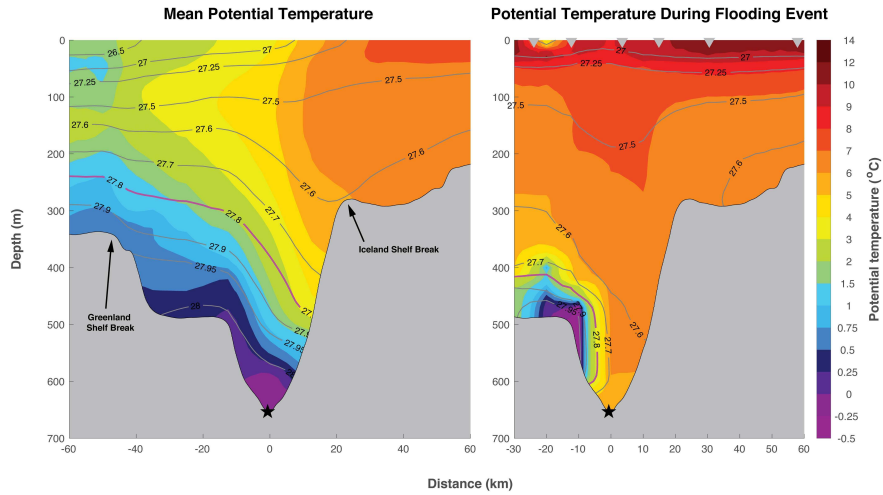
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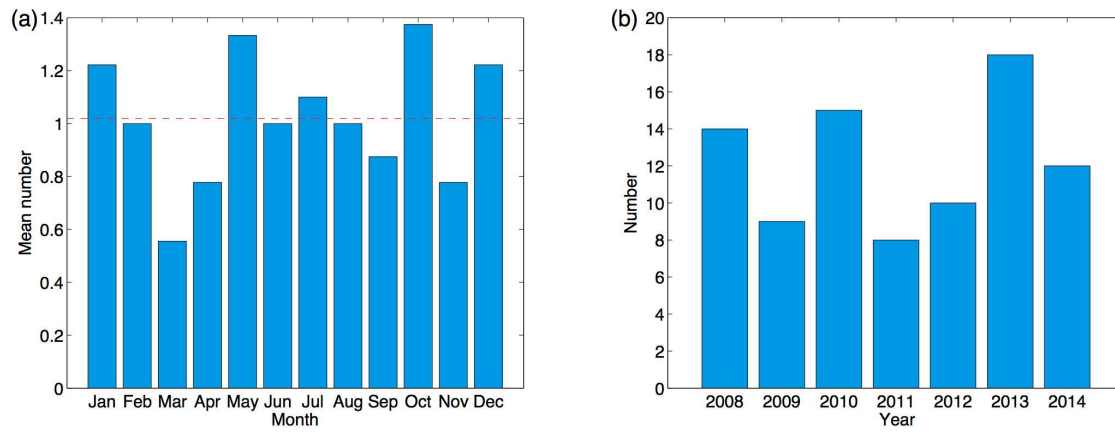
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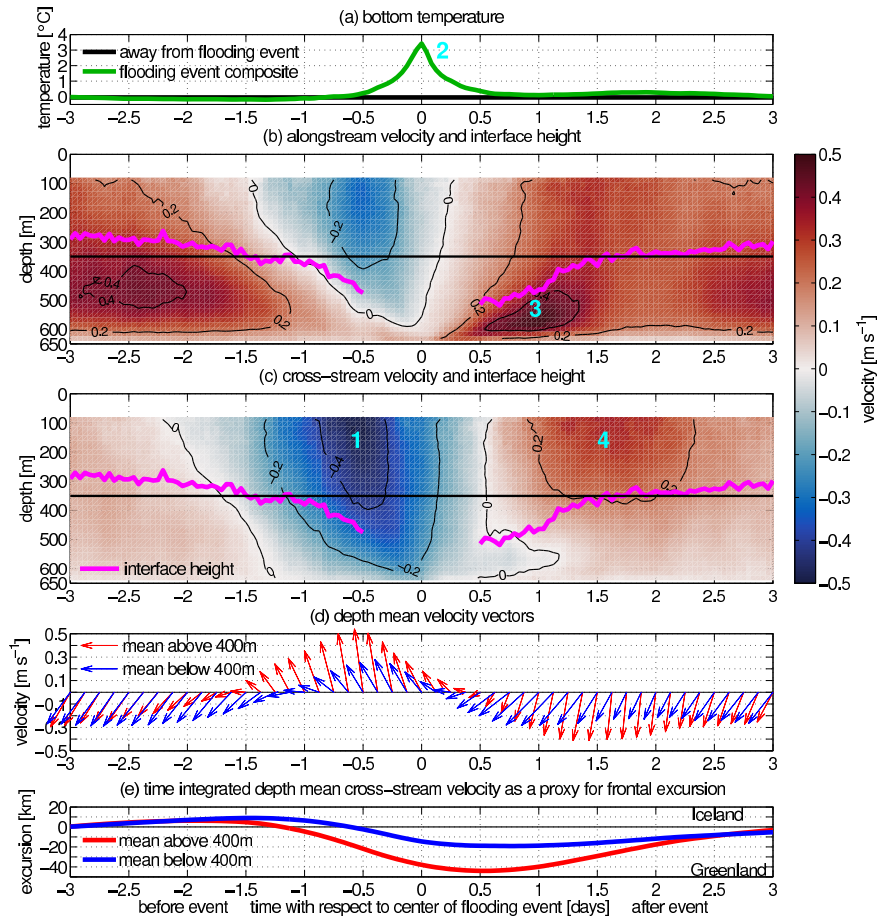
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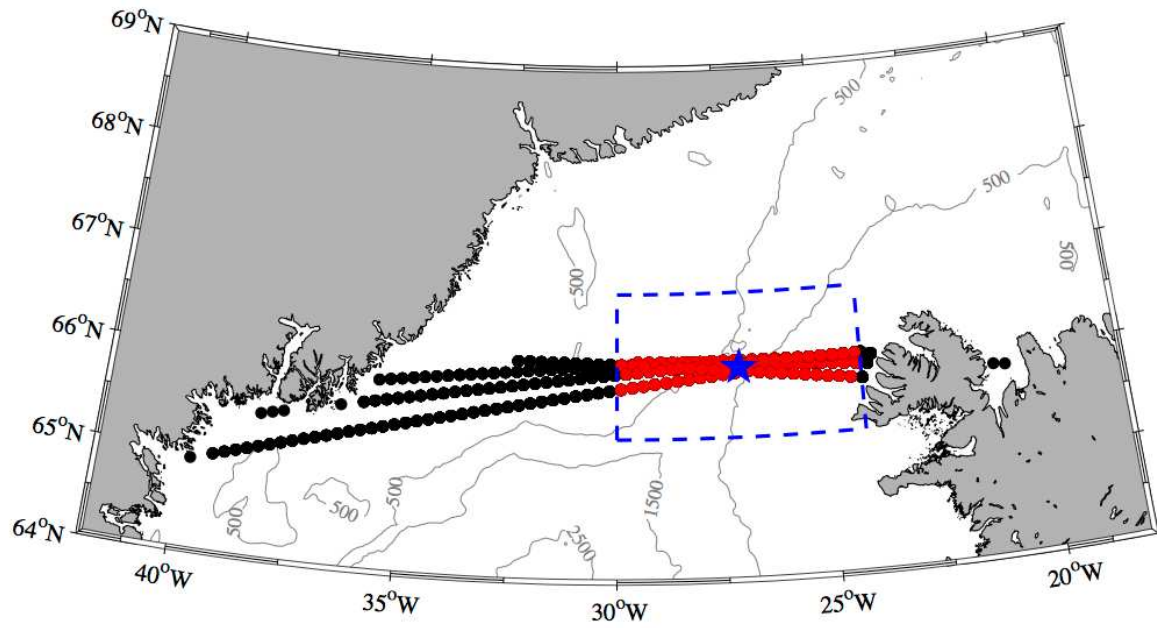
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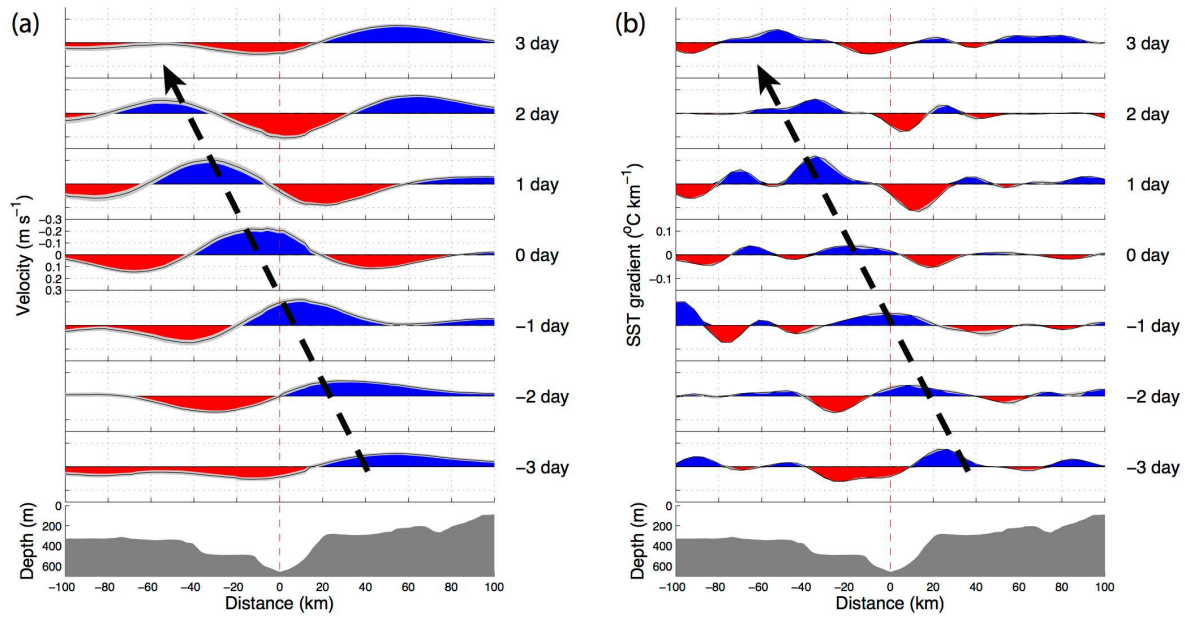
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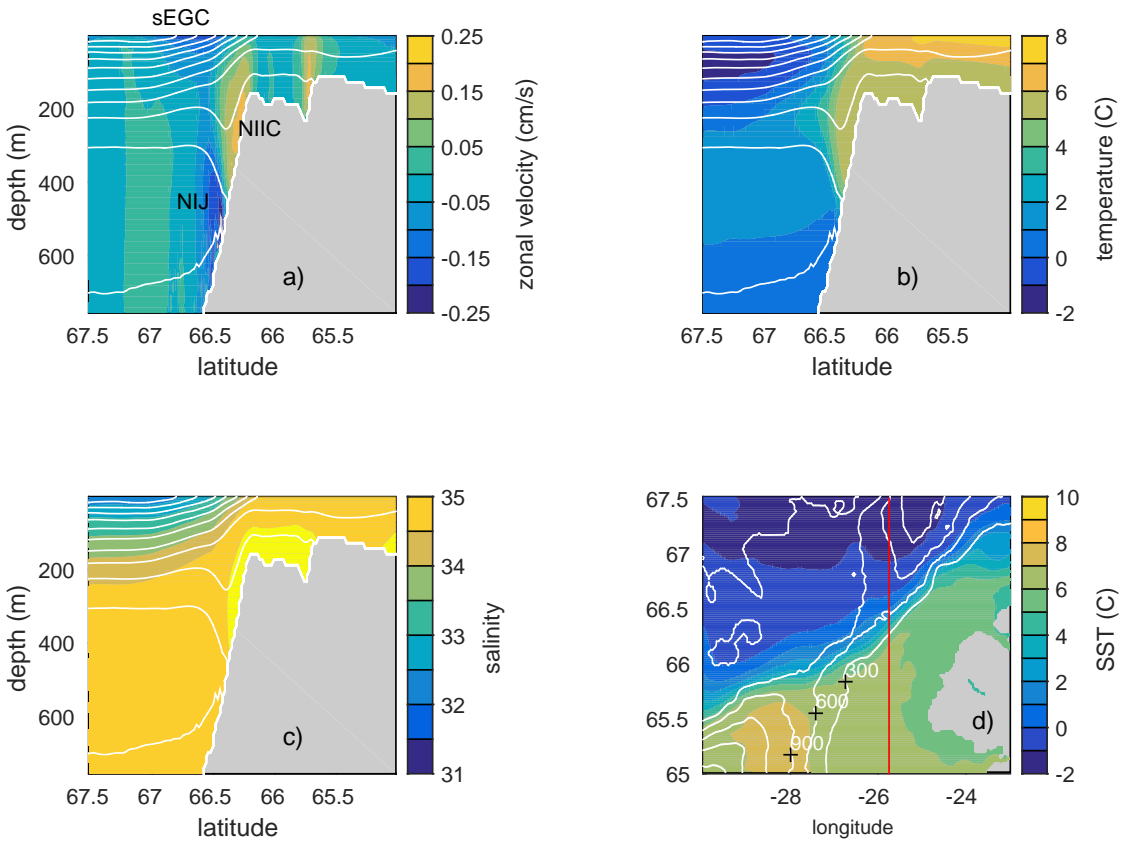
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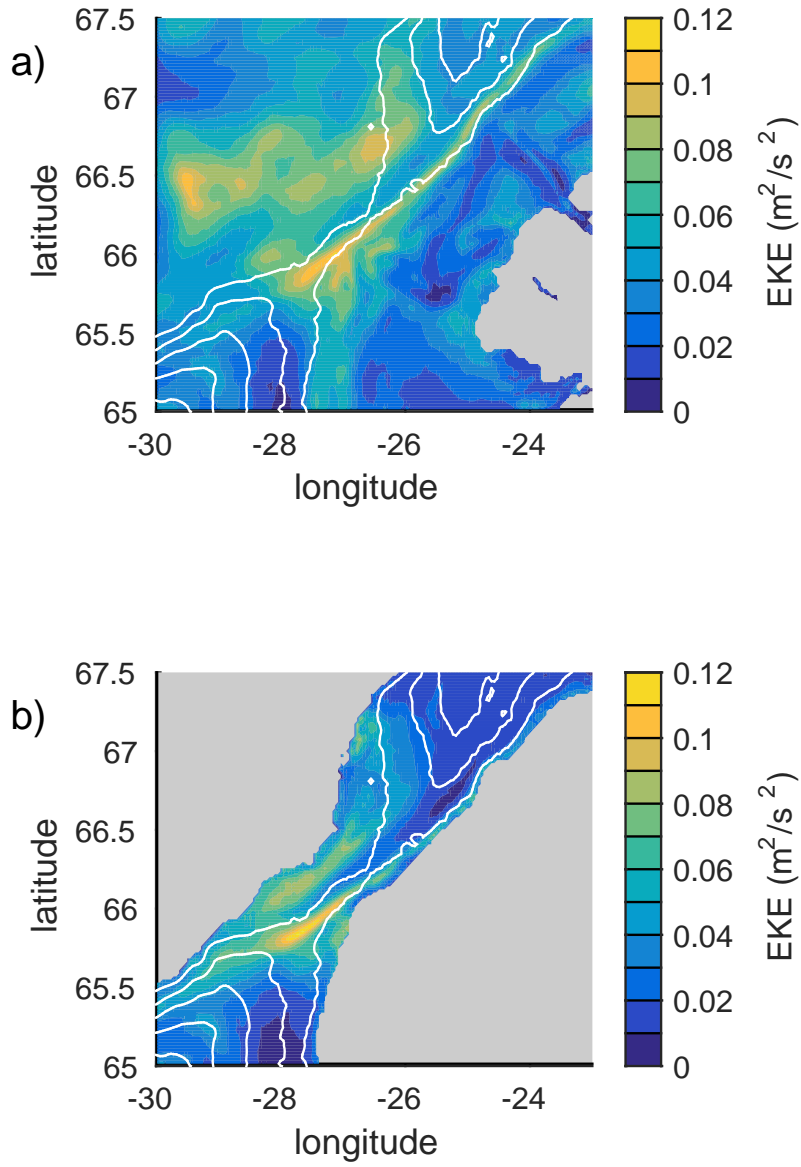
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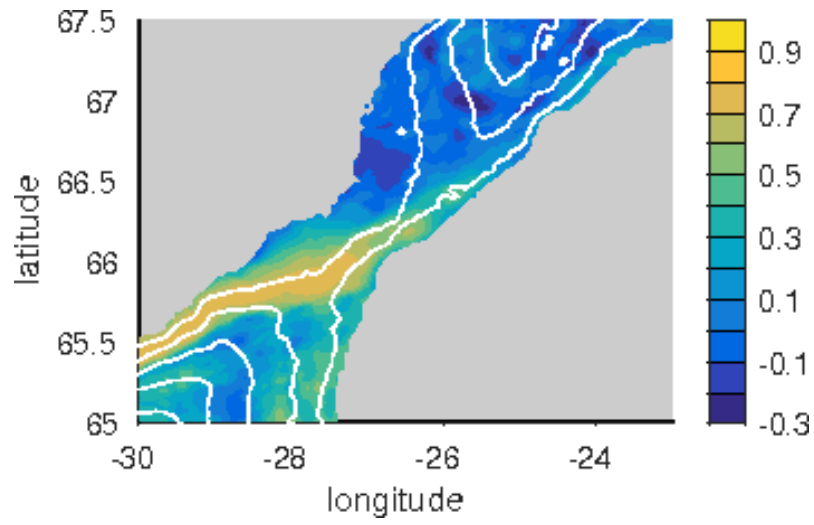
638 FIG. 6. (a) Composite along-strait surface geostrophic velocity during a flooding event. Blue is poleward and
 639 red is equatorward. Day 0 corresponds to the peak of the event, defined using the DS1 mooring data. The dashed
 640 black arrow denotes the westward propagation of the NIIC. The bottom panel shows the cross-strait bathymetry.
 641 The red dashed line marks the location of the DS1 mooring. (b) Composite of the cross-strait SST gradient
 642 during a flooding event. The positive gradient associated with the NIIC (large blue peak) propagates westward
 643 at the same speed as the surface geostrophic velocity signal in (a), indicated by the identical dashed black arrow.



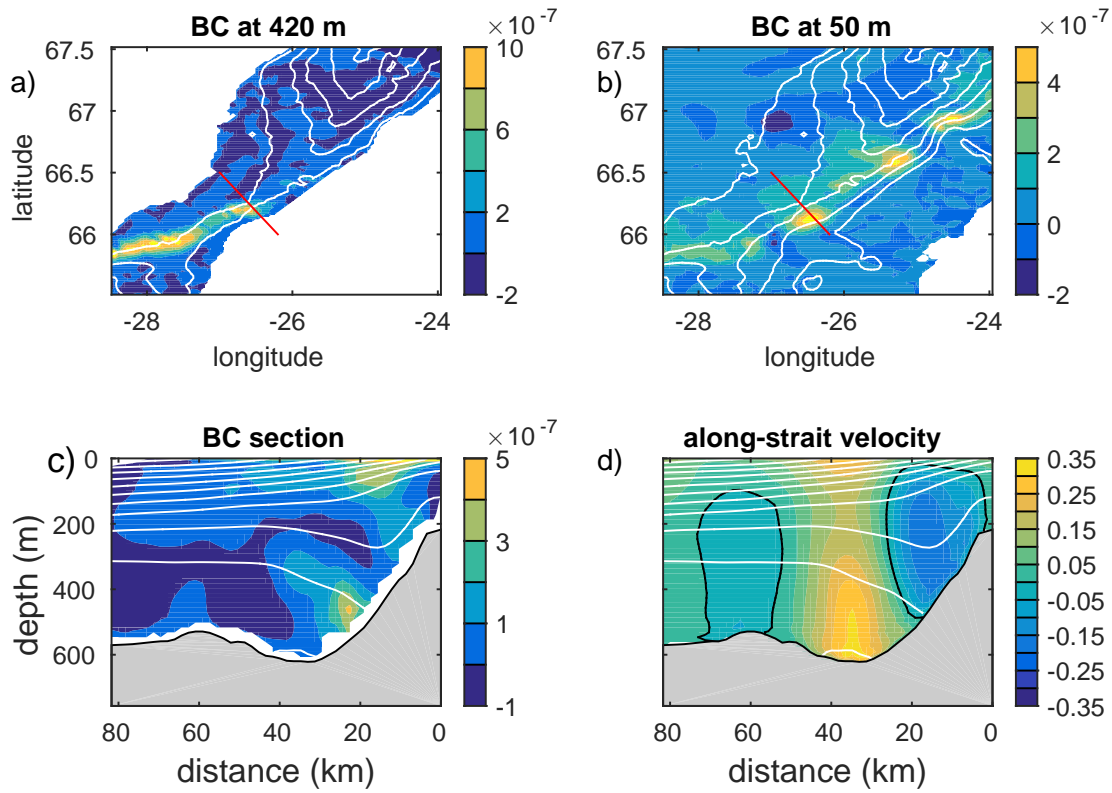
644 FIG. 7. Meridional section of the mean a) zonal velocity ($m s^{-1}$); b) temperature (C); c) salinity. The white
 645 contours are potential density, contour interval $0.2 kg m^{-3}$. d) mean sea surface temperature. The white contours
 646 are bottom topography (contour interval 300 m), and the red line marks the location of the meridional sections
 647 in a-c.



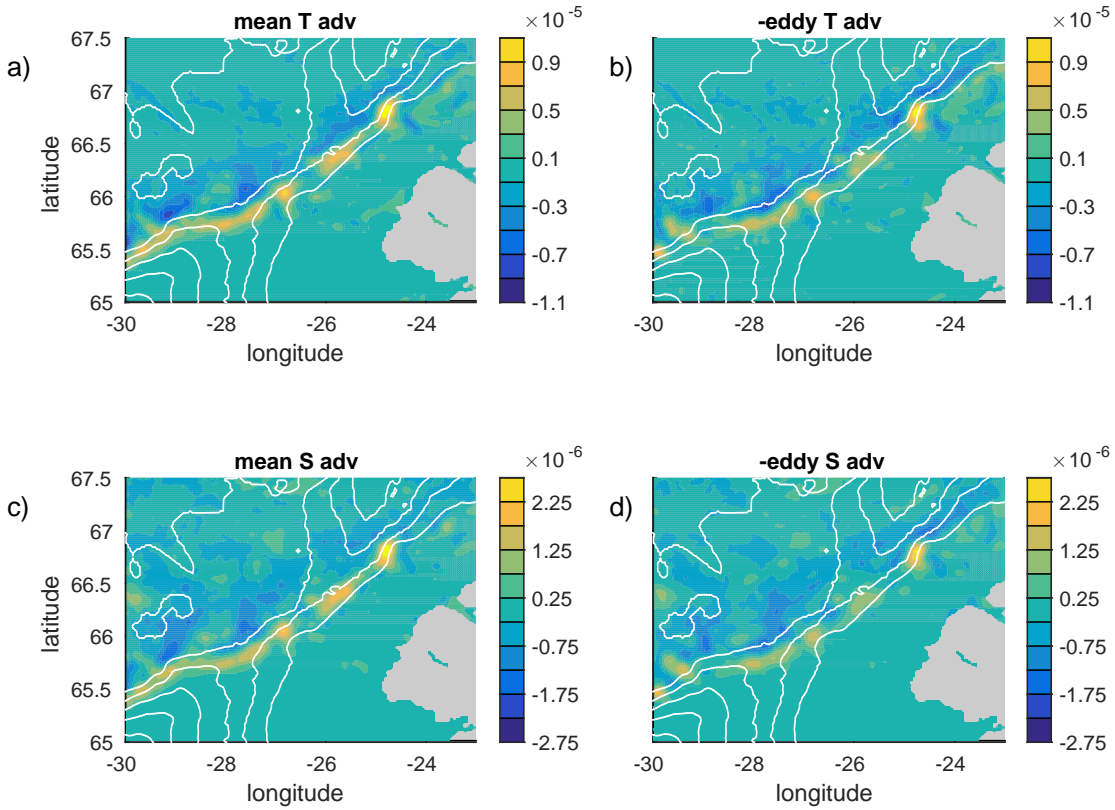
648 FIG. 8. Eddy kinetic energy at a) 50 m and b) 420 m depth ($m^2 s^{-2}$). Bottom topography is indicated by the
 649 white contours (contour interval 300 m).



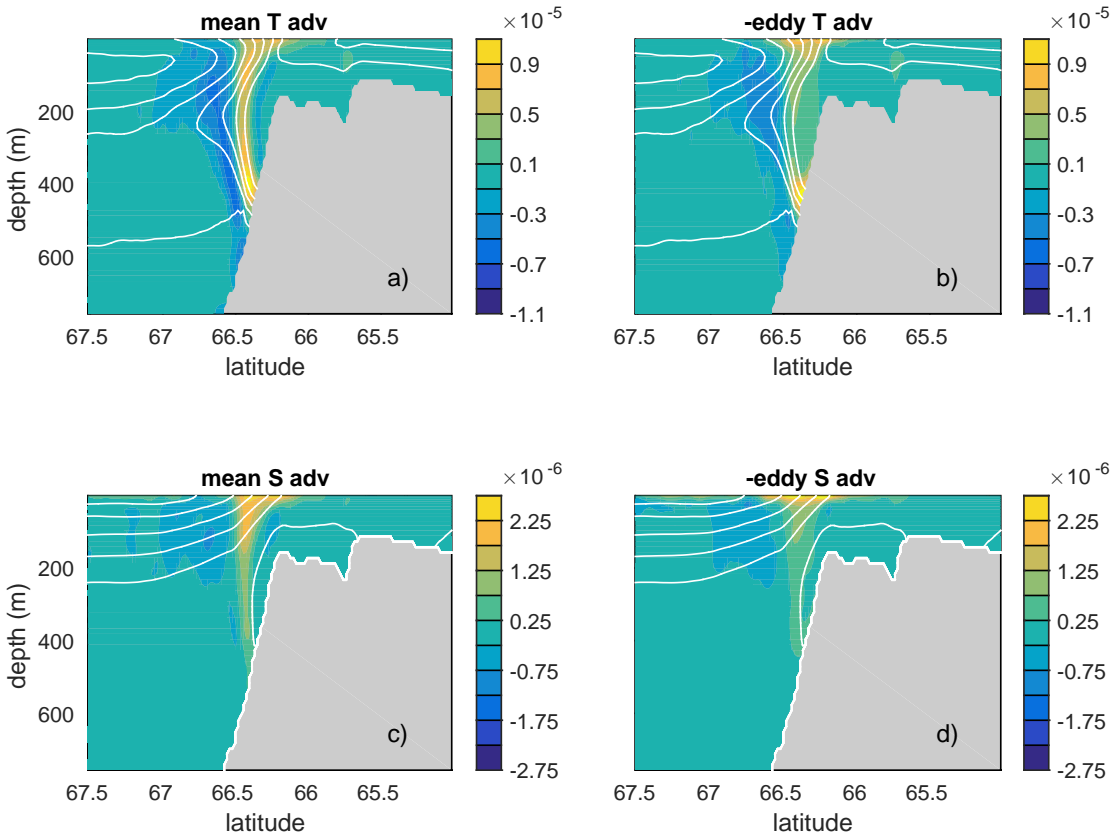
650 FIG. 9. Correlation between the temperature anomaly at 50 m depth and the temperature anomaly at 420 m
651 depth. Bottom topography is indicated by the white contours (contour interval 300 m)



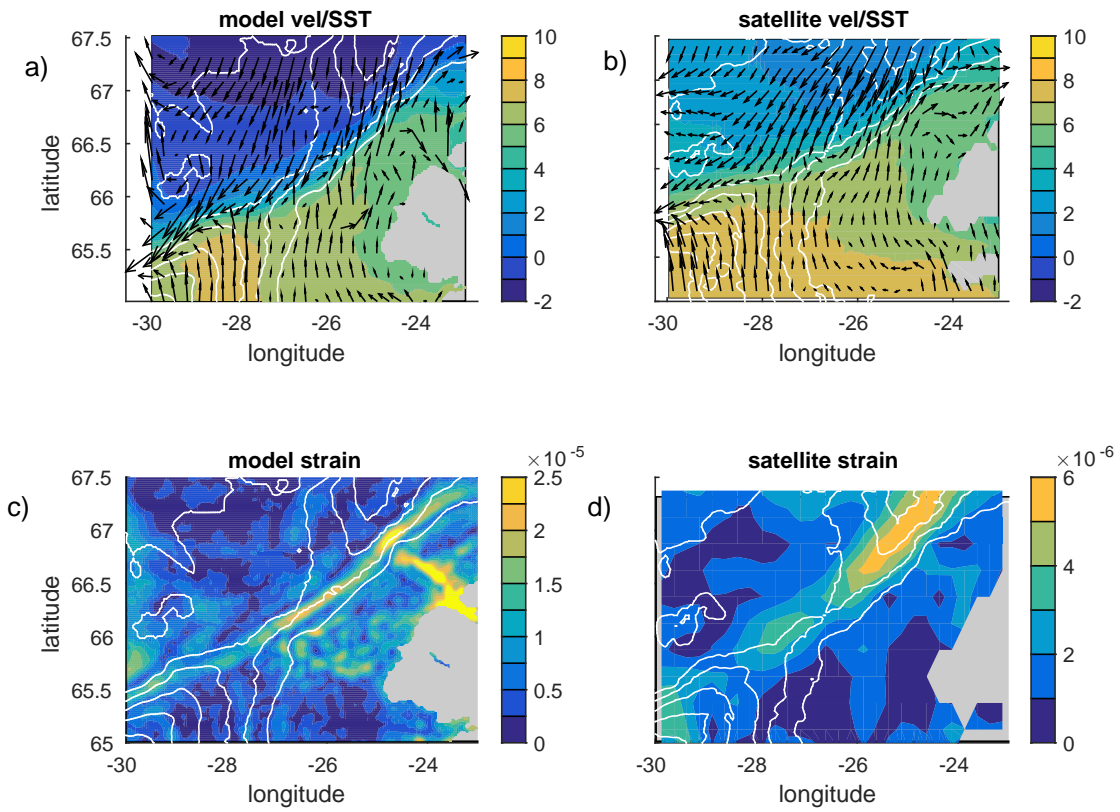
652 FIG. 10. Baroclinic energy conversion rate BC at a) 420 m, b) 50m ($m^2 s^{-3}$). c) vertical section of BC along
 653 the red line in a and b. d) time-mean velocity normal to the section ($m s^{-1}$, zero contour is black; distance is in
 654 km from southeast to northwest; positive is flow towards the southwest).



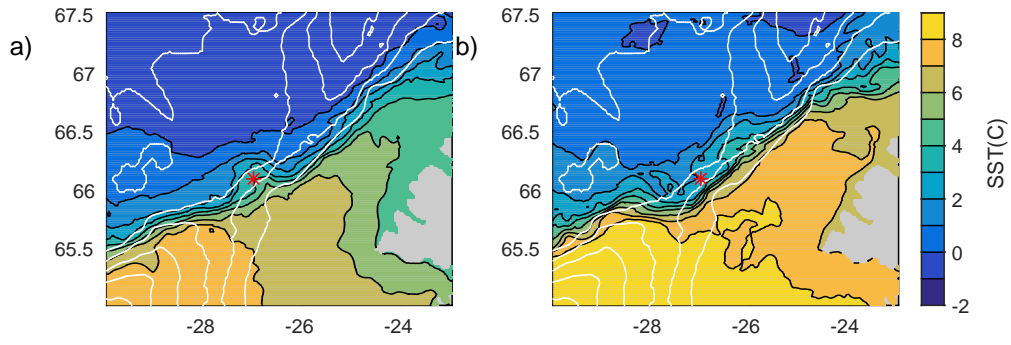
655 FIG. 11. Contributions to the temperature ($C s^{-1}$) and salinity (s^{-1}) tendency equations at 50 m depth due to
 656 a) mean temperature advection; b) negative of the eddy temperature advection; c) mean salinity advection; d)
 657 negative of the eddy salinity advection.



658 FIG. 12. Contributions to the temperature and salinity tendency equations across the meridional section at
 659 $25.75^{\circ}W$ (see Fig. 7d) due to a) mean temperature advection; b) negative of the eddy temperature advection; c)
 660 mean salinity advection; d) negative of the eddy salinity advection. The white contours are for temperature in a)
 661 and b) (contour interval $1^{\circ}C$) and salinity in c) and d) (contour interval 0.5).



662 FIG. 13. a) Mean velocity and SST from the model; b) same as a) calculated using the satellite data; c) mean
 663 strain field from the model (s^{-1}); d) same as c) using the satellite derived surface velocity.



664 FIG. 14. Composite averages of SST from the model for (a) periods when flooding events were present, and
 665 (b) periods when DSOW boluses were present. White contours are bottom topography (contour interval 300 m).
 666 The red asterisk is the location of the mooring.