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Definition of functionals of the geopotential used in GrafLab software

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Gravitational potential

$$V(r, \theta, \lambda) = \frac{GM}{r} \left[1 + \sum_{n=2}^M \left(\frac{R}{r} \right)^n \sum_{m=0}^n (\bar{C}_{n,m} \cos(m\lambda) + \bar{S}_{n,m} \sin(m\lambda)) \bar{P}_{n,m}(\cos \theta) \right] \quad (1)$$

Gravitational tensor in the spherical coordinates

$$\mathbf{V}(r, \theta, \lambda) = \begin{pmatrix} V_{rr} & V_{r\theta} & V_{r\lambda} \\ V_{\theta r} & V_{\theta\theta} & V_{\theta\lambda} \\ V_{\lambda r} & V_{\lambda\theta} & V_{\lambda\lambda} \end{pmatrix} \quad (2)$$

$$\begin{aligned} V_{rr}(r, \theta, \lambda) &= \frac{\partial^2 V(r, \theta, \lambda)}{\partial r^2} \\ &= \frac{GM}{r^3} \left[2 + \sum_{n=2}^M \left(\frac{R}{r} \right)^n (n+1)(n+2) \sum_{m=0}^n (\bar{C}_{n,m} \cos(m\lambda) + \bar{S}_{n,m} \sin(m\lambda)) \bar{P}_{n,m}(\cos \theta) \right] \end{aligned} \quad (3)$$

$$\begin{aligned} V_{r\theta}(r, \theta, \lambda) &= \frac{1}{r} \frac{\partial^2 V(r, \theta, \lambda)}{\partial r \partial \theta} \\ &= -\frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r} \right)^n (n+1) \sum_{m=0}^n (\bar{C}_{n,m} \cos(m\lambda) + \bar{S}_{n,m} \sin(m\lambda)) \frac{d\bar{P}_{n,m}(\cos \theta)}{d\theta} \end{aligned} \quad (4)$$

$$\begin{aligned} V_{r\lambda}(r, \theta, \lambda) &= \frac{1}{r \sin \theta} \frac{\partial^2 V(r, \theta, \lambda)}{\partial r \partial \lambda} \\ &= -\frac{GM}{r^3 \sin \theta} \sum_{n=2}^M \left(\frac{R}{r} \right)^n (n+1) \sum_{m=0}^n (\bar{S}_{n,m} \cos(m\lambda) - \bar{C}_{n,m} \sin(m\lambda)) m \bar{P}_{n,m}(\cos \theta) \end{aligned} \quad (5)$$

$$\begin{aligned} V_{\theta\theta}(r, \theta, \lambda) &= \frac{1}{r^2} \frac{\partial^2 V(r, \theta, \lambda)}{\partial \theta^2} \\ &= \frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r} \right)^n \sum_{m=0}^n (\bar{C}_{n,m} \cos(m\lambda) + \bar{S}_{n,m} \sin(m\lambda)) \frac{d^2 \bar{P}_{n,m}(\cos \theta)}{d\theta^2} \end{aligned} \quad (6)$$

$$\begin{aligned} V_{\theta\lambda}(r, \theta, \lambda) &= \frac{1}{r^2 \sin \theta} \frac{\partial^2 V(r, \theta, \lambda)}{\partial \theta \partial \lambda} \\ &= \frac{GM}{r^3 \sin \theta} \sum_{n=2}^M \left(\frac{R}{r} \right)^n \sum_{m=0}^n (\bar{S}_{n,m} \cos(m\lambda) - \bar{C}_{n,m} \sin(m\lambda)) m \frac{d\bar{P}_{n,m}(\cos \theta)}{d\theta} \end{aligned} \quad (7)$$

$$\begin{aligned}
V_{\lambda\lambda}(r, \theta, \lambda) &= \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 V(r, \theta, \lambda)}{\partial \lambda^2} \\
&= -\frac{GM}{r^3 \sin^2 \theta} \sum_{n=2}^M \left(\frac{R}{r}\right)^n \sum_{m=0}^n (\bar{C}_{n,m} \cos(m\lambda) + \bar{S}_{n,m} \sin(m\lambda)) m^2 \bar{P}_{n,m}(\cos \theta)
\end{aligned} \tag{8}$$

Gravitational tensor in the local north-oriented reference frame¹

$$\mathbf{V}(r, \theta, \lambda) = \begin{pmatrix} V_{xx} & V_{xy} & V_{xz} \\ V_{yx} & V_{yy} & V_{yz} \\ V_{zx} & V_{zy} & V_{zz} \end{pmatrix} \tag{9}$$

$$\begin{aligned}
V_{xx} = & -\frac{GM}{r^3} + \frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r}\right)^n \sum_{m=-n}^n \bar{C}_{n,m} Q_m(\lambda) \left(a_{n,m} \bar{P}_{n,|m|-2}(\cos \theta) \right. \\
& \left. + [b_{n,m} - (n+1)(n+2)] \bar{P}_{n,|m|}(\cos \theta) + c_{n,m} \bar{P}_{n,|m|+2}(\cos \theta) \right)
\end{aligned} \tag{10}$$

$$\begin{aligned}
V_{yy} = & -\frac{GM}{r^3} - \frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r}\right)^n \sum_{m=-n}^n \bar{C}_{n,m} Q_m(\lambda) \left(a_{n,m} \bar{P}_{n,|m|-2}(\cos \theta) \right. \\
& \left. + b_{n,m} \bar{P}_{n,|m|}(\cos \theta) + c_{n,m} \bar{P}_{n,|m|+2}(\cos \theta) \right)
\end{aligned} \tag{11}$$

$$\begin{aligned}
V_{xy} = & \frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r}\right)^n \sum_{m=-n}^n \bar{C}_{n,m} Q_{-m}(\lambda) \left(d_{n,m} \bar{P}_{n-1,|m|-2}(\cos \theta) \right. \\
& \left. + g_{n,m} \bar{P}_{n-1,|m|}(\cos \theta) + h_{n,m} \bar{P}_{n-1,|m|+2}(\cos \theta) \right), \quad m \neq 0
\end{aligned} \tag{12}$$

$$\begin{aligned}
V_{xz} = & \frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r}\right)^n \sum_{m=-n}^n \bar{C}_{n,m} Q_m(\lambda) \left(\beta_{n,m} \bar{P}_{n,|m|-1}(\cos \theta) \right. \\
& \left. + \gamma_{n,m} \bar{P}_{n,|m|+1}(\cos \theta) \right)
\end{aligned} \tag{13}$$

$$\begin{aligned}
V_{yz} = & \frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r}\right)^n \sum_{m=-n}^n \bar{C}_{n,m} Q_{-m}(\lambda) \left(\mu_{n,m} \bar{P}_{n-1,|m|-1}(\cos \theta) \right. \\
& \left. + \nu_{n,m} \bar{P}_{n-1,|m|+1}(\cos \theta) \right), \quad m \neq 0
\end{aligned} \tag{14}$$

$$V_{zz} = 2\frac{GM}{r^3} + \frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r}\right)^n (n+1)(n+2) \sum_{m=-n}^n \bar{C}_{n,m} Q_m(\lambda) \bar{P}_{n,|m|}(\cos \theta) \tag{15}$$

where

$$Q_m(\lambda) = \begin{cases} \cos m\lambda, & m \geq 0 \\ \sin |m|\lambda, & m < 0 \end{cases} \tag{16}$$

¹In the GrafLab, Eq. (10) - Eq. (15) have been slightly modified, see appendix A.

$$a_{n,m} = 0, \quad |m| = 0, 1 \quad (17)$$

$$a_{n,m} = \frac{\sqrt{1 + \delta_{|m|,2}}}{4} \sqrt{n^2 - (|m| - 1)^2} \sqrt{n + |m|} \sqrt{n - |m| + 2}, \quad 2 \leq |m| \leq n \quad (18)$$

$$b_{n,m} = \frac{(n + |m| + 1)(n + |m| + 2)}{2(|m| + 1)}, \quad |m| = 0, 1 \quad (19)$$

$$b_{n,m} = \frac{n^2 + m^2 + 3n + 2}{2}, \quad 2 \leq |m| \leq n \quad (20)$$

$$c_{n,m} = \frac{\sqrt{1 + \delta_{|m|,0}}}{4} \sqrt{n^2 - (|m| + 1)^2} \sqrt{n - |m|} \sqrt{n + |m| + 2}, \quad |m| = 0, 1 \quad (21)$$

$$c_{n,m} = \frac{1}{4} \sqrt{n^2 - (|m| + 1)^2} \sqrt{n - |m|} \sqrt{n + |m| + 2}, \quad 2 \leq |m| \leq n \quad (22)$$

$$d_{n,m} = 0, \quad |m| = 1 \quad (23)$$

$$d_{n,m} = -\frac{m}{4|m|} \sqrt{\frac{2n+1}{2n-1}} \sqrt{1 + \delta_{|m|,2}} \sqrt{n^2 - (|m| - 1)^2} \\ \times \sqrt{n + |m|} \sqrt{n + |m| - 2}, \quad 2 \leq |m| \leq n \quad (24)$$

$$g_{n,m} = \frac{m}{4|m|} \sqrt{\frac{2n+1}{2n-1}} \sqrt{n+1} \sqrt{n-1} (n+2), \quad |m| = 1 \quad (25)$$

$$g_{n,m} = \frac{m}{2} \sqrt{\frac{2n+1}{2n-1}} \sqrt{n + |m|} \sqrt{n - |m|}, \quad 2 \leq |m| \leq n \quad (26)$$

$$h_{n,m} = \frac{m}{4|m|} \sqrt{\frac{2n+1}{2n-1}} \sqrt{n-3} \sqrt{n-2} \sqrt{n-1} \sqrt{n+2}, \quad |m| = 1 \quad (27)$$

$$h_{n,m} = \frac{m}{4|m|} \sqrt{\frac{2n+1}{2n-1}} \sqrt{n^2 - (|m| + 1)^2} \sqrt{n - |m|} \sqrt{n - |m| - 2}, \quad 2 \leq |m| \leq n \quad (28)$$

$$\beta_{n,m} = 0, \quad m = 0 \quad (29)$$

$$\beta_{n,m} = \frac{n+2}{2} \sqrt{1 + \delta_{|m|,1}} \sqrt{n + |m|} \sqrt{n - |m| + 1}, \quad 1 \leq |m| \leq n \quad (30)$$

$$\gamma_{n,m} = -(n+2) \sqrt{\frac{n(n+1)}{2}}, \quad m = 0 \quad (31)$$

$$\gamma_{n,m} = -\frac{n+2}{2} \sqrt{n - |m|} \sqrt{n + |m| + 1}, \quad 1 \leq |m| \leq n \quad (32)$$

$$\mu_{n,m} = -\frac{m}{|m|} \left(\frac{n+2}{2} \right) \sqrt{\frac{2n+1}{2n-1}} \sqrt{1 + \delta_{|m|,1}} \sqrt{n + |m|} \sqrt{n + |m| - 1} \quad (33)$$

$$\nu_{n,m} = -\frac{m}{|m|} \left(\frac{n+2}{2} \right) \sqrt{\frac{2n+1}{2n-1}} \sqrt{n - |m|} \sqrt{n - |m| - 1} \quad (34)$$

$$\delta_{p,q} = \begin{cases} 1, & p = q, \\ 0, & p \neq q. \end{cases} \quad (35)$$

Gravity potential

$$\begin{aligned}
 W(r, \theta, \lambda) = & \frac{GM}{r} \left[1 + \sum_{n=2}^M \left(\frac{R}{r} \right)^n \sum_{m=0}^n (\overline{C}_{n,m} \cos(m\lambda) + \overline{S}_{n,m} \sin(m\lambda)) \overline{P}_{n,m}(\cos \theta) \right] \\
 & + \frac{1}{2} \omega^2 r^2 \sin^2 \theta
 \end{aligned} \tag{36}$$

Gravity

$$\begin{aligned}
 g(r, \theta, \lambda) = & |\nabla W(r, \theta, \lambda)| \\
 = & \sqrt{\left(\frac{\partial V}{\partial r} + \frac{\partial V_c}{\partial r} \right)^2 + \left[\frac{1}{r} \left(\frac{\partial V}{\partial \theta} + \frac{\partial V_c}{\partial \theta} \right) \right]^2 + \left[\frac{1}{r \sin \theta} \left(\frac{\partial V}{\partial \lambda} + \frac{\partial V_c}{\partial \lambda} \right) \right]^2},
 \end{aligned} \tag{37}$$

$$\begin{aligned}
 \frac{\partial V}{\partial r} = & -\frac{GM}{r^2} \left[1 + \sum_{n=2}^M \left(\frac{R}{r} \right)^n (n+1) \sum_{m=0}^n (\overline{C}_{n,m} \cos(m\lambda) + \overline{S}_{n,m} \sin(m\lambda)) \overline{P}_{n,m}(\cos \theta) \right] \\
 \frac{\partial V}{\partial \theta} = & \frac{GM}{r} \sum_{n=2}^M \left(\frac{R}{r} \right)^n \sum_{m=0}^n (\overline{C}_{n,m} \cos(m\lambda) + \overline{S}_{n,m} \sin(m\lambda)) \frac{d\overline{P}_{n,m}(\cos \theta)}{d\theta} \\
 \frac{\partial V}{\partial \lambda} = & \frac{GM}{r} \sum_{n=2}^M \left(\frac{R}{r} \right)^n \sum_{m=0}^n (\overline{S}_{n,m} \cos(m\lambda) - \overline{C}_{n,m} \sin(m\lambda)) m \overline{P}_{n,m}(\cos \theta)
 \end{aligned} \tag{38}$$

$$V_c = \frac{1}{2} \omega^2 r^2 \sin^2 \theta \tag{39}$$

$$\frac{\partial V_c}{\partial r} = \omega^2 r \sin^2 \theta, \quad \frac{\partial V_c}{\partial \theta} = \omega^2 r^2 \sin \theta \cos \theta, \quad \frac{\partial V_c}{\partial \lambda} = 0 \tag{40}$$

Gravity sa (spherical approximation)

$$g_{sa}(r, \theta, \lambda) = \sqrt{\left(\frac{\partial V}{\partial r} + \frac{\partial V_c}{\partial r} \right)^2} \tag{41}$$

Second radial derivative of gravity potential

$$\frac{\partial^2 W(r, \theta, \lambda)}{\partial r^2} = \frac{\partial^2 V(r, \theta, \lambda)}{\partial r^2} + \frac{\partial^2 V_c(r, \theta, \lambda)}{\partial r^2} \quad (42)$$

$$\begin{aligned} & \frac{\partial^2 V(r, \theta, \lambda)}{\partial r^2} \\ &= \frac{GM}{r^3} \left[2 + \sum_{n=2}^M \left(\frac{R}{r} \right)^n (n+1)(n+2) \sum_{m=0}^n (\bar{C}_{n,m} \cos(m\lambda) + \bar{S}_{n,m} \sin(m\lambda)) \bar{P}_{n,m}(\cos \theta) \right] \end{aligned} \quad (43)$$

$$\frac{\partial^2 V_c(r, \theta, \lambda)}{\partial r^2} = \omega^2 \sin^2 \theta \quad (44)$$

Disturbing potential

$$T(r, \theta, \lambda) = \frac{GM}{r} \sum_{n=2}^M \left(\frac{R}{r} \right)^n \sum_{m=0}^n (\Delta \bar{C}_{n,m} \cos(m\lambda) + \Delta \bar{S}_{n,m} \sin(m\lambda)) \bar{P}_{n,m}(\cos \theta) \quad (45)$$

$$\Delta \bar{C}_{n,m} = \bar{C}_{n,m} - \bar{C}_{n,m}^{Ell} \frac{GM^{Ell}}{GM} \left(\frac{a^{Ell}}{R} \right)^n \quad (46)$$

$$\Delta \bar{S}_{n,m} = \bar{S}_{n,m} - \bar{S}_{n,m}^{Ell} = \bar{S}_{n,m} \quad (47)$$

Gravity disturbance

$$\delta g(r, \theta, \lambda) = g(r, \theta, \lambda) - \gamma_{SH}(r, \theta) \quad (48)$$

in which

- $\gamma_{SH}(r, \theta, \lambda)$ is the normal gravity evaluated from the spherical harmonics. The same formulae as for $g(r, \theta, \lambda)$ hold for computing $\gamma_{SH}(r, \theta, \lambda)$ by replacing $\bar{C}_{n,m}$, $\bar{S}_{n,m}$ by $\bar{C}_{n,m}^{Ell}$, $\bar{S}_{n,m}^{Ell}$ in eq. (37) – (38).

Gravity disturbance sa (spherical approximation)

$$\begin{aligned} \delta g_{sa}(r, \theta, \lambda) &= - \frac{\partial T(r, \theta, \lambda)}{\partial r} \\ &= \frac{GM}{r^2} \sum_{n=2}^M \left(\frac{R}{r} \right)^n (n+1) \sum_{m=0}^n (\Delta \bar{C}_{n,m} \cos(m\lambda) + \Delta \bar{S}_{n,m} \sin(m\lambda)) \bar{P}_{n,m}(\cos \theta) \end{aligned} \quad (49)$$

Gravity anomaly sa (spherical approximation)

$$\begin{aligned}
\Delta g_{sa}(r, \theta, \lambda) &= -\frac{\partial T(r, \theta, \lambda)}{\partial r} - \frac{2}{r}T(r, \theta, \lambda) \\
&= \frac{GM}{r^2} \sum_{n=2}^M \left(\frac{R}{r}\right)^n (n-1) \sum_{m=0}^n (\Delta \bar{C}_{n,m} \cos(m\lambda) + \Delta \bar{S}_{n,m} \sin(m\lambda)) \bar{P}_{n,m}(\cos \theta)
\end{aligned} \tag{50}$$

Second radial derivative of disturbing potential

$$\begin{aligned}
T_{rr}(r, \theta, \lambda) &= \frac{\partial^2 T(r, \theta, \lambda)}{\partial r^2} \\
&= \frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r}\right)^n (n+1)(n+2) \sum_{m=0}^n (\Delta \bar{C}_{n,m} \cos(m\lambda) + \Delta \bar{S}_{n,m} \sin(m\lambda)) \bar{P}_{n,m}(\cos \theta)
\end{aligned} \tag{51}$$

Disturbing tensor in the spherical coordinates

$$\mathbf{T}(r, \theta, \lambda) = \begin{pmatrix} T_{rr} & T_{r\theta} & T_{r\lambda} \\ T_{\theta r} & T_{\theta\theta} & T_{\theta\lambda} \\ T_{\lambda r} & T_{\lambda\theta} & T_{\lambda\lambda} \end{pmatrix} \tag{52}$$

$$\begin{aligned}
T_{rr}(r, \theta, \lambda) &= \frac{\partial^2 T(r, \theta, \lambda)}{\partial r^2} \\
&= \frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r}\right)^n (n+1)(n+2) \sum_{m=0}^n (\Delta \bar{C}_{n,m} \cos(m\lambda) + \Delta \bar{S}_{n,m} \sin(m\lambda)) \bar{P}_{n,m}(\cos \theta)
\end{aligned} \tag{53}$$

$$\begin{aligned}
T_{r\theta}(r, \theta, \lambda) &= \frac{1}{r} \frac{\partial^2 T(r, \theta, \lambda)}{\partial r \partial \theta} \\
&= -\frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r}\right)^n (n+1) \sum_{m=0}^n (\Delta \bar{C}_{n,m} \cos(m\lambda) + \Delta \bar{S}_{n,m} \sin(m\lambda)) \frac{d\bar{P}_{n,m}(\cos \theta)}{d\theta}
\end{aligned} \tag{54}$$

$$\begin{aligned}
T_{r\lambda}(r, \theta, \lambda) &= \frac{1}{r \sin \theta} \frac{\partial^2 T(r, \theta, \lambda)}{\partial r \partial \lambda} \\
&= -\frac{GM}{r^3 \sin \theta} \sum_{n=2}^M \left(\frac{R}{r}\right)^n (n+1) \sum_{m=0}^n (\Delta \bar{S}_{n,m} \cos(m\lambda) - \Delta \bar{C}_{n,m} \sin(m\lambda)) m \bar{P}_{n,m}(\cos \theta)
\end{aligned} \tag{55}$$

$$\begin{aligned}
T_{\theta\theta}(r, \theta, \lambda) &= \frac{1}{r^2} \frac{\partial^2 T(r, \theta, \lambda)}{\partial \theta^2} \\
&= \frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r}\right)^n \sum_{m=0}^n (\Delta \bar{C}_{n,m} \cos(m\lambda) + \Delta \bar{S}_{n,m} \sin(m\lambda)) \frac{d^2 \bar{P}_{n,m}(\cos \theta)}{d\theta^2}
\end{aligned} \tag{56}$$

$$\begin{aligned}
T_{\theta\lambda}(r, \theta, \lambda) &= \frac{1}{r^2 \sin \theta} \frac{\partial^2 T(r, \theta, \lambda)}{\partial \theta \partial \lambda} \\
&= \frac{GM}{r^3 \sin \theta} \sum_{n=2}^M \left(\frac{R}{r}\right)^n \sum_{m=0}^n (\Delta \bar{S}_{n,m} \cos(m\lambda) - \Delta \bar{C}_{n,m} \sin(m\lambda)) m \frac{d\bar{P}_{n,m}(\cos \theta)}{d\theta}
\end{aligned} \tag{57}$$

$$\begin{aligned}
T_{\lambda\lambda}(r, \theta, \lambda) &= \frac{1}{r^2 \sin^2 \theta} \frac{\partial^2 T(r, \theta, \lambda)}{\partial \lambda^2} \\
&= -\frac{GM}{r^3 \sin^2 \theta} \sum_{n=2}^M \left(\frac{R}{r}\right)^n \sum_{m=0}^n (\Delta \bar{C}_{n,m} \cos(m\lambda) + \Delta \bar{S}_{n,m} \sin(m\lambda)) m^2 \bar{P}_{n,m}(\cos \theta)
\end{aligned} \tag{58}$$

Disturbing tensor in the local north-oriented reference frame²

$$\mathbf{T}(r, \theta, \lambda) = \begin{pmatrix} T_{xx} & T_{xy} & T_{xz} \\ T_{yx} & T_{yy} & T_{yz} \\ T_{zx} & T_{zy} & T_{zz} \end{pmatrix} \tag{59}$$

$$\begin{aligned}
T_{xx} &= \frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r}\right)^n \sum_{m=-n}^n \bar{C}_{n,m} Q_m(\lambda) \left(a_{n,m} \bar{P}_{n,|m|-2}(\cos \theta) \right. \\
&\quad \left. + [b_{n,m} - (n+1)(n+2)] \bar{P}_{n,|m|}(\cos \theta) + c_{n,m} \bar{P}_{n,|m|+2}(\cos \theta) \right)
\end{aligned} \tag{60}$$

²In the GrafLab, Eq. (60) - Eq. (65) have been slightly modified, see appendix A.

$$T_{yy} = -\frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r}\right)^n \sum_{m=-n}^n \bar{C}_{n,m} Q_m(\lambda) \left(a_{n,m} \bar{P}_{n,|m|-2}(\cos \theta) \right. \\ \left. + b_{n,m} \bar{P}_{n,|m|}(\cos \theta) + c_{n,m} \bar{P}_{n,|m|+2}(\cos \theta) \right) \quad (61)$$

$$T_{xy} = \frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r}\right)^n \sum_{m=-n}^n \bar{C}_{n,m} Q_{-m}(\lambda) \left(d_{n,m} \bar{P}_{n-1,|m|-2}(\cos \theta) \right. \\ \left. + g_{n,m} \bar{P}_{n-1,|m|}(\cos \theta) + h_{n,m} \bar{P}_{n-1,|m|+2}(\cos \theta) \right), \quad m \neq 0 \quad (62)$$

$$T_{xz} = \frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r}\right)^n \sum_{m=-n}^n \bar{C}_{n,m} Q_m(\lambda) \left(\beta_{n,m} \bar{P}_{n,|m|-1}(\cos \theta) \right. \\ \left. + \gamma_{n,m} \bar{P}_{n,|m|+1}(\cos \theta) \right) \quad (63)$$

$$T_{yz} = \frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r}\right)^n \sum_{m=-n}^n \bar{C}_{n,m} Q_{-m}(\lambda) \left(\mu_{n,m} \bar{P}_{n-1,|m|-1}(\cos \theta) \right. \\ \left. + \nu_{n,m} \bar{P}_{n-1,|m|+1}(\cos \theta) \right), \quad m \neq 0 \quad (64)$$

$$T_{zz} = \frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r}\right)^n (n+1)(n+2) \sum_{m=-n}^n \bar{C}_{n,m} Q_m(\lambda) \bar{P}_{n,|m|}(\cos \theta) \quad (65)$$

Deflections of the vertical

$$\xi(r, \theta, \lambda) = -\frac{GM}{r^2 \gamma(r, \theta)} \sum_{n=2}^M \left(\frac{R}{r}\right)^n \sum_{m=0}^n (\Delta \bar{C}_{n,m} \cos(m\lambda) \\ + \Delta \bar{S}_{n,m} \sin(m\lambda)) \frac{d\bar{P}_{n,m}(\cos \theta)}{d\theta} \quad (66)$$

$$\eta(r, \theta, \lambda) = -\frac{GM}{r^2 \gamma(r, \theta) \sin \theta} \sum_{n=2}^M \left(\frac{R}{r}\right)^n \sum_{m=0}^n (\Delta \bar{S}_{n,m} \cos(m\lambda) \\ - \Delta \bar{C}_{n,m} \sin(m\lambda)) m \bar{P}_{n,m}(\cos \theta) \quad (67)$$

$$\Theta(r, \theta, \lambda) = \sqrt{\xi^2(r, \theta, \lambda) + \eta^2(r, \theta, \lambda)} \quad (68)$$

Geoid undulation

$$H(\theta, \lambda) = \sum_{n=0}^M \sum_{m=0}^n (\overline{HC}_{n,m} \cos(m\lambda) + \overline{HS}_{n,m} \sin(m\lambda)) \bar{P}_{n,m}(\cos \theta) \quad (69)$$

$$N(\theta, \lambda) = \frac{T(r_{ell}, \theta, \lambda) - 2\pi G \rho H^2(\theta, \lambda)}{\gamma(r_{ell}, \theta)} \quad (70)$$

where

- G denotes the Newtonian gravitational constant,
 $G = 6.67259 \cdot 10^{-11} \text{ m}^3 \cdot \text{kg}^{-1} \cdot \text{s}^{-2}$
- ρ denotes the density of the crust, $\rho = 2670 \text{ kg} \cdot \text{m}^{-3}$
- $r_{ell} = r_{ell}(\theta)$ denotes the radius-coordinate of point Q_0 on the reference ellipsoid

Height anomaly ell

$$\zeta_{ell}(r, \theta, \lambda) = \frac{T(r, \theta, \lambda)}{\gamma(r, \theta)} \quad (71)$$

Height anomaly

$$\zeta(r, \theta, \lambda) = \zeta_{ell}(r_{ell}, \theta, \lambda) - \delta g_{sa}(r_{ell}, \theta, \lambda) \frac{H(\theta, \lambda) + N(\theta, \lambda)}{\gamma(r_{ell}, \theta)} \quad (72)$$

References

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- Petrovskaya, M.S., Vershkov, A.N., 2006. Non-singular expressions for the gravity gradients in the local north-oriented and orbital reference frames. Journal of Geodesy 80, 117-127. doi: 10.1029/2011JB008916.

A Modified non-singular expressions for the gravity gradients in the LNOF

In this appendix, we will demonstrate how we modified the non-singular expressions for the gravity gradients in the LNOF (Eq. (10) - Eq. (15) and Eq. (60) - Eq. (65)). As an example, let us mentioned only one particular expression for the element T_{xx} , which has the following form

$$\begin{aligned}
T_{xx} = & \frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r} \right)^n \sum_{m=0}^n (\bar{C}_{nm} \cos m\lambda + \bar{S}_{nm} \sin m\lambda) \\
& \times \left(a_{nm} \bar{P}_{n,m-2}(\cos \theta) \right. \\
& + [b_{nm} - (n+1)(n+2)] \bar{P}_{nm}(\cos \theta) \\
& \left. + c_{nm} \bar{P}_{n,m+2}(\cos \theta) \right),
\end{aligned} \tag{73}$$

in which

$$a_{nm} = 0, \quad m = 0, 1 \tag{74}$$

$$\begin{aligned}
a_{nm} = & \frac{\sqrt{1 + \delta_{m,2}}}{4} \sqrt{n^2 - (m-1)^2} \\
& \times \sqrt{n+m} \sqrt{n-m+2}, \quad 2 \leq m \leq n
\end{aligned} \tag{75}$$

$$b_{nm} = \frac{(n+m+1)(n+m+2)}{2(m+1)}, \quad m = 0, 1 \tag{76}$$

$$b_{nm} = \frac{n^2 + m^2 + 3n + 2}{2}, \quad 2 \leq m \leq n \tag{77}$$

$$\begin{aligned}
c_{nm} = & \frac{\sqrt{1 + \delta_{m,0}}}{4} \sqrt{n^2 - (m+1)^2} \sqrt{n-m} \\
& \times \sqrt{n+m+2}, \quad m = 0, 1
\end{aligned} \tag{78}$$

$$c_{nm} = \frac{1}{4} \sqrt{n^2 - (m+1)^2} \sqrt{n-m} \sqrt{n+m+2}, \quad 2 \leq m \leq n \tag{79}$$

$$\delta_{p,q} = \begin{cases} 1, & p = q, \\ 0, & p \neq q. \end{cases} \tag{80}$$

From Eq. (73), one can see that in addition to the term $\bar{P}_{nm}(\cos \theta)$, two other terms $\bar{P}_{n,m-2}(\cos \theta)$ and $\bar{P}_{n,m+2}(\cos \theta)$ must be computed for each m . From the practical numerical point of view, this is not an issue, if fixed-degree recursions have been used to evaluate fnALFs. In this case, for each m these terms have been already computed with the term $\bar{P}_{nm}(\cos \theta)$ essentially. In the GrafLab, however, we used fixed-order recursions, which are more frequently used in geodesy. In this case, with every change of m in the order-dependent loop, it is necessary to evaluate not only the term $\bar{P}_{nm}(\cos \theta)$, but also the two other terms. In other words, redundant

computations occur. Thus, we modified Eq. (73) in the way that only the term $\bar{P}_{nm}(\cos \theta)$ is needed to be computed. We present Eq. (73) in the following form

$$\begin{aligned}
T_{xx} = & \frac{GM}{r^3} \sum_{n=2}^M \left(\frac{R}{r} \right)^n \\
& \times \sum_{m=0}^n \left[\left(\bar{C}_{n,m+2} \cos(m+2)\lambda + \bar{S}_{n,m+2} \sin(m+2)\lambda \right) a_{n,m+2} \right. \\
& + \left(\bar{C}_{nm} \cos m\lambda + \bar{S}_{nm} \sin m\lambda \right) (b_{nm} - (n+1)(n+2)) \\
& \left. + \left(\bar{C}_{n,m-2} \cos(m-2)\lambda + \bar{S}_{n,m-2} \sin(m-2)\lambda \right) c_{n,m-2} \right] \\
& \times \bar{P}_{nm}(\cos \theta),
\end{aligned} \tag{81}$$

where

$$\left. \begin{array}{l} \bar{C}_{n,m+2} \\ \bar{S}_{n,m+2} \\ \cos(m+2)\lambda \\ \sin(m+2)\lambda \\ a_{n,m+2} \end{array} \right\} = 0, \quad m+2 > n, \tag{82}$$

$$\left. \begin{array}{l} \bar{C}_{n,m-2} \\ \bar{S}_{n,m-2} \\ \cos(m-2)\lambda \\ \sin(m-2)\lambda \\ c_{n,m-2} \end{array} \right\} = 0, \quad m-2 < 0. \tag{83}$$

The main idea of Eq. (81) is that the set of spherical harmonic coefficients is usually stored during the whole computational process, hence the coefficients $\bar{C}_{n,m+2}, \bar{S}_{n,m+2}$ and $\bar{C}_{n,m-2}, \bar{S}_{n,m-2}$ may be simply restored when necessary instead of the redundant computation of $\bar{P}_{n,m-2}(\cos \theta)$ and $\bar{P}_{n,m+2}(\cos \theta)$ in Eq. (73). The formulae for the remaining elements $T_{yy}, T_{zz}, T_{xy}, T_{xz}, T_{yz}$ may be easily modified in the same way.